

MRI Systems III: Gradients

M219 - Principles and Applications of MRI

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1/24/2022

Course Overview

- Course website
 - <https://mrri.ucla.edu/pages/m219>
- Course schedule
 - https://mrri.ucla.edu/pages/m219_2022
- Assignments
 - Homework #1 due on 1/26 by 5pm
 - Homework #2 will be out on 1/26

Course Overview

- Office Hours

- TA (Ran Yan) - Tuesday 4-5pm

[https://uclahs.zoom.us/j/96870184581?](https://uclahs.zoom.us/j/96870184581?pwd=VkcZL0lyRkxsQ3FHcnIxQ1M2U3hPdZ09)

[pwd=VkcZL0lyRkxsQ3FHcnIxQ1M2U3hPdZ09](https://uclahs.zoom.us/j/96870184581?pwd=VkcZL0lyRkxsQ3FHcnIxQ1M2U3hPdZ09)

Password: 900645

- Instructor (Kyung Sung) - Friday 2-3pm

[https://uclahs.zoom.us/j/94058312815?](https://uclahs.zoom.us/j/94058312815?pwd=Tkl3ajhkamdGTnhqOVNnbk5RMnJGQT09)

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Password: 888767

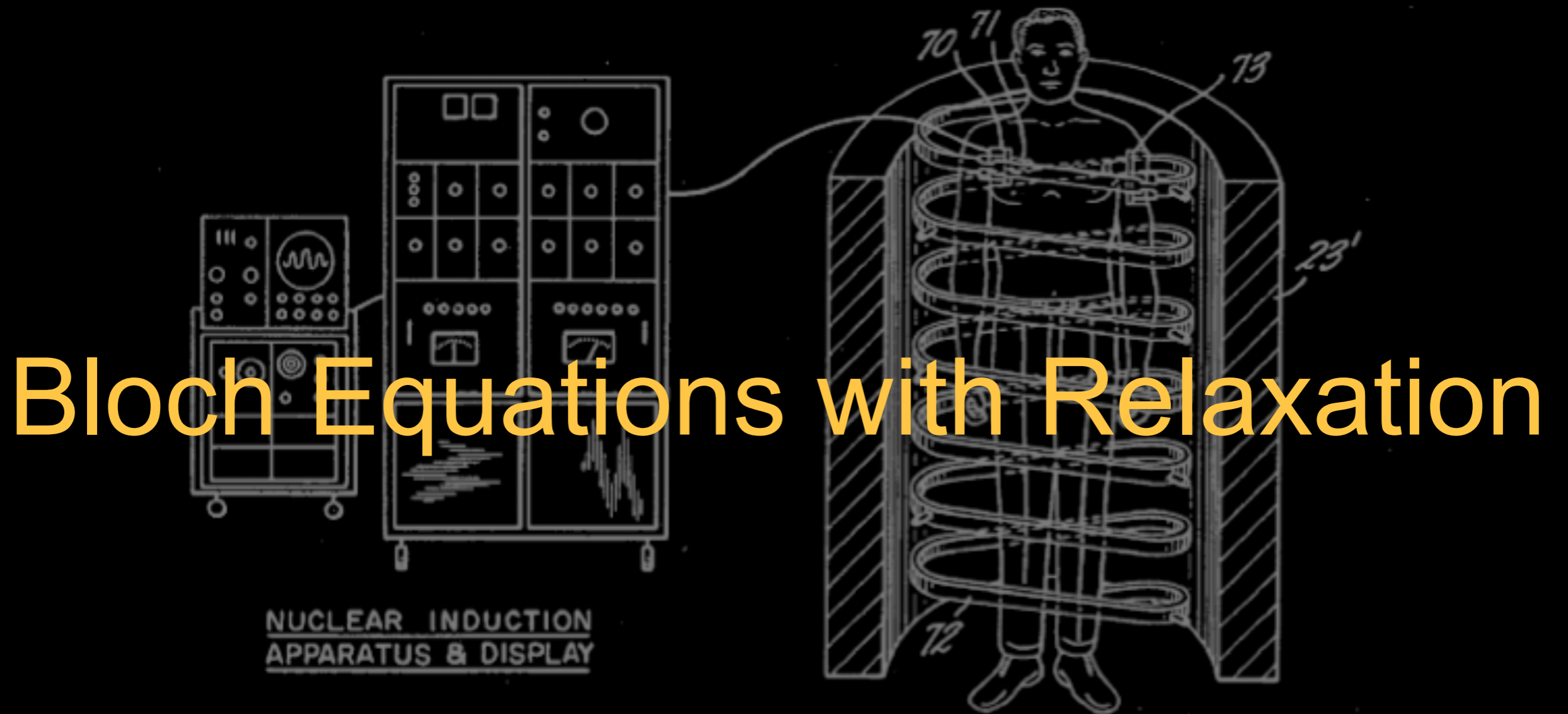


FIG. 2

Bloch Equations with Relaxation



Bloch Equations with Relaxation

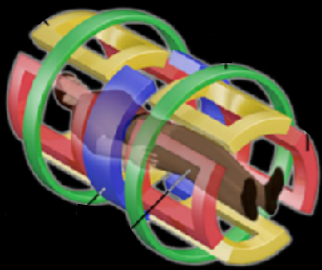
$$\frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B} - \frac{M_x \hat{i} + M_y \hat{j}}{T_2} - \frac{(M_z - M_0) \hat{k}}{T_1}$$

- **Differential Equation**
 - Ordinary, Coupled, Non-linear
- **No analytic solution, in general.**
 - Analytic solutions for simple cases.
 - Numerical solutions for all cases.
- **Phenomenological**
 - Exponential behavior is an approximation.

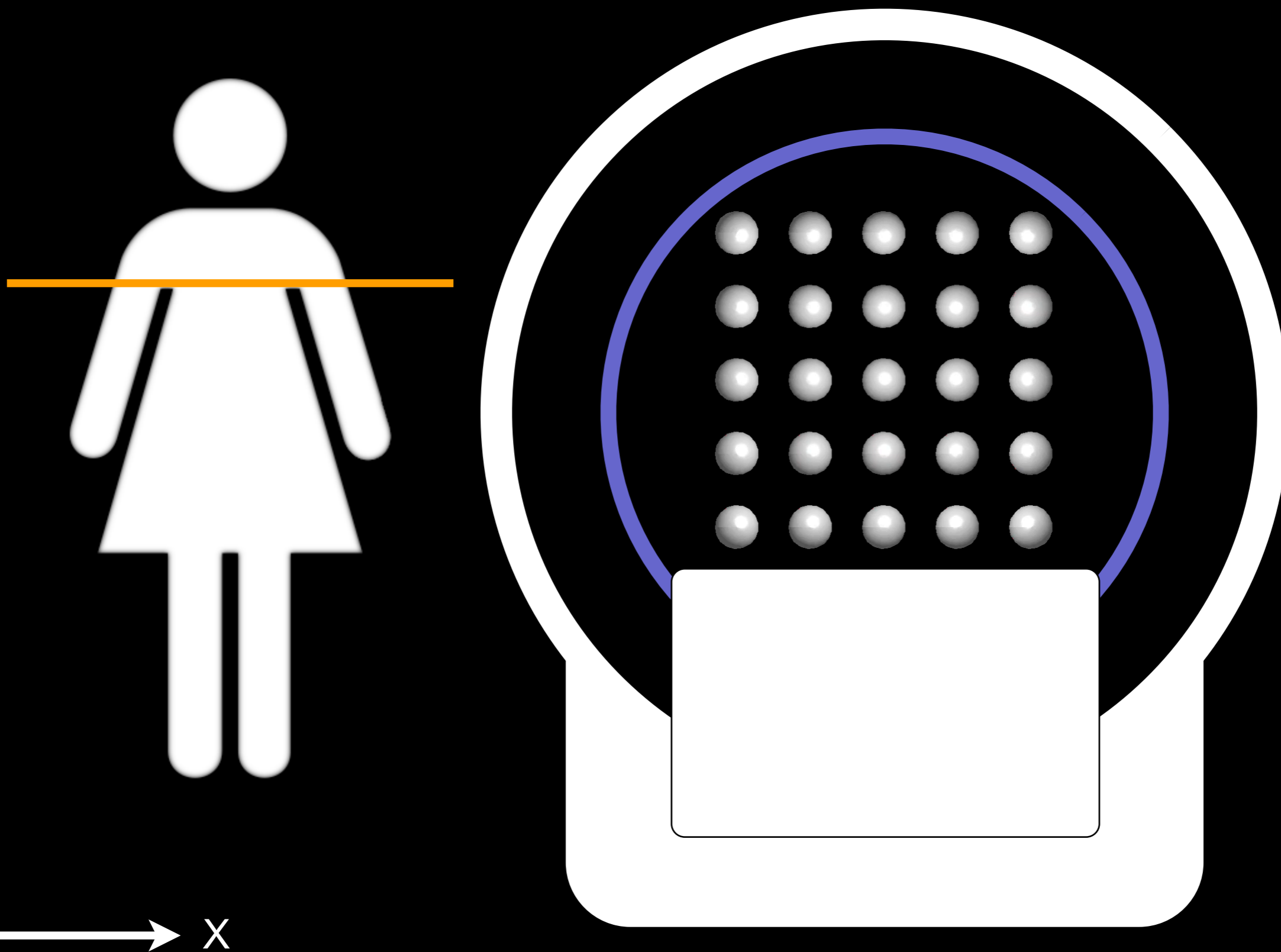
Bloch Equations - Lab Frame

$$\frac{d\vec{M}}{dt} = \underbrace{\vec{M} \times \gamma \vec{B}}_{\text{Precession}} - \underbrace{\frac{M_x \hat{i} + M_y \hat{j}}{T_2}}_{\text{Transverse Relaxation}} - \underbrace{\frac{(M_z - M_0) \hat{k}}{T_1}}_{\text{Longitudinal Relaxation}}$$

- **Precession**
 - Magnitude of M unchanged
 - Phase (rotation) of M changes due to B
- **Relaxation**
 - T_1 changes are slow O(100ms)
 - T_2 changes are fast O(10ms)
 - Magnitude of M can be ZERO
- **Diffusion**
 - Spins are thermodynamically driven to exchange positions.
 - Bloch-Torrey Equations



Excitation and Relaxation



The magnetization relaxes after excitation (forced precession).

Bloch Equations – Rotating Frame

$$\frac{\partial \vec{M}_{rot}}{\partial t} = \underbrace{\gamma \vec{M}_{rot} \times \vec{B}_{eff}}_{\text{“Precession”}} - \underbrace{\frac{M_{x'} \vec{i}' + M_{y'} \vec{j}'}{T_2}}_{\text{Transverse Relaxation}} - \underbrace{\frac{(M_{z'} - M_0) \vec{k}'}{T_1}}_{\text{Longitudinal Relaxation}}$$

$$\vec{B}_{eff} \triangleq \frac{\vec{\omega}}{\gamma} + \vec{B}_{rot}$$

Effective B-field that M experiences in the rotating frame

The applied B₀ and B₁ field in the rotating frame

Fictitious field created by the rotating frame that demodulates the apparent effect of B₀

Free Precession in the Rotating Frame with Relaxation

Free Precession in the Rotating Frame

$$\frac{\partial \vec{M}_{rot}}{\partial t} = \gamma \vec{M}_{rot} \times \vec{B}_{eff} - \frac{M_{x'} \vec{i}' + M_{y'} \vec{j}'}{T_2} - \frac{(M_{z'} - M_0) \vec{k}'}{T_1}$$

$$\vec{B}_{eff} \triangleq \frac{\vec{\omega}}{\gamma} + \vec{B}_{rot}$$

$$\vec{\omega}_{rot} = \vec{\omega} = -\gamma B_0 \hat{k} \quad \vec{B}_{rot} = B_0 \hat{k}$$

$$\vec{B}_{eff} = \vec{0}$$

$$\frac{\partial \vec{M}_{rot}}{\partial t} = - \frac{M_{x'} \vec{i}' + M_{y'} \vec{j}'}{T_2} - \frac{(M_{z'} - M_0) \vec{k}'}{T_1}$$

Free Precession in the Rotating Frame

$$\frac{\partial \vec{M}_{rot}}{\partial t} = \underbrace{-\frac{M_{x'}\vec{i}' + M_{y'}\vec{j}'}{T_2}}_{\text{Transverse Relaxation}} - \underbrace{\frac{(M_{z'} - M_0)\vec{k}'}{T_1}}_{\text{Longitudinal Relaxation}}$$

- **No precession**
- **T₁ and T₂ Relaxation**
- **Drop the diffusion term**
- **System of first order, linear, separable ODEs!**

Free Precession in the Rotating Frame

$$\frac{\partial \vec{M}_{rot}}{\partial t} = \underbrace{-\frac{M_{x'}\vec{i}' + M_{y'}\vec{j}'}{T_2}}_{\text{Transverse Relaxation}} - \underbrace{\frac{(M_{z'} - M_0)\vec{k}'}{T_1}}_{\text{Longitudinal Relaxation}}$$

Solution:

$$M_{z'}(t) = M_z^0 e^{-t/T_1} + M_0 (1 - e^{-t/T_1})$$

$$M_{x'y'}(t) = M_{x'y'}(0_+) e^{-t/T_2}$$

Forced Precession in the Rotating Frame with Relaxation

Forced Precession in the Rot. Frame with Relaxation

$$\frac{\partial \vec{M}_{rot}}{\partial t} = \gamma \vec{M}_{rot} \times \vec{B}_{eff} - \frac{M_{x'} \vec{i}' + M_{y'} \vec{j}'}{T_2} - \frac{(M_{z'} - M_0) \vec{k}'}{T_1}$$

$$\vec{B}_{eff} \triangleq \frac{\vec{\omega}}{\gamma} + \vec{B}_{rot}$$

$$\vec{\omega}_{rot} = \vec{\omega} = -\gamma B_0 \hat{k} \quad \vec{B}_{rot} = B_0 \hat{k} + B_1^e(t) \hat{i}'$$

$$\vec{B}_{eff} = B_1^e(t) \hat{i}'$$

Forced Precession in the Rot. Frame with Relaxation

$$\frac{\partial \vec{M}_{rot}}{\partial t} = \gamma \vec{M}_{rot} \times \vec{B}_{eff} - \frac{M_{x'} \vec{i}' + M_{y'} \vec{j}'}{T_2} - \frac{(M_{z'} - M_0) \vec{k}'}{T_1}$$

$$\vec{B}_{eff} = B_1^e(t) \hat{i}'$$

- **B1 induced nutation**
- **T₁ and T₂ Relaxation**
- **Drop the diffusion term**
- **System or first order, linear, coupled PDEs!**
- **When does this equation apply?**

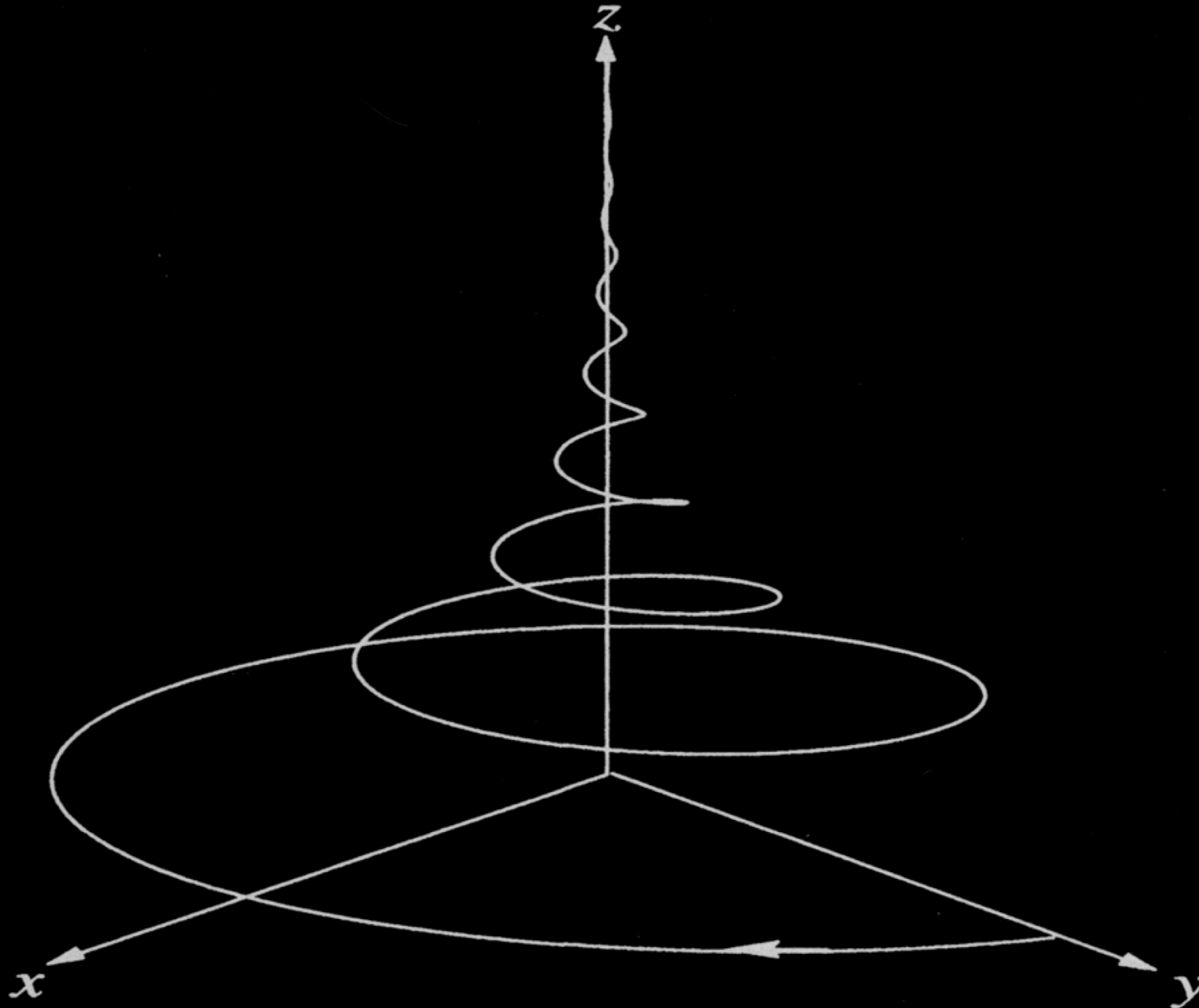
Forced Precession in the Rotating Frame with Relaxation

- RF pulses are short
 - $100\mu\text{s}$ to 5ms
- Relaxation time constants are long
 - T_1 $O(100\text{s})$ ms
 - T_2 $O(10\text{s})$ ms
- Complicated Coupling
- Best suited for simulation

Free? Forced? Relaxation?

- **We've considered all combinations of:**
 - **Free and forced precession**
 - **With and without relaxation**
 - **Laboratory and rotating frames**
- **Which one's concern M219 the most?**
 - **Free precession in the rotating frame with relaxation**
 - **Forced precession in the rotating frame without relaxation.**
- **We can, in fact, simulate all of them...**

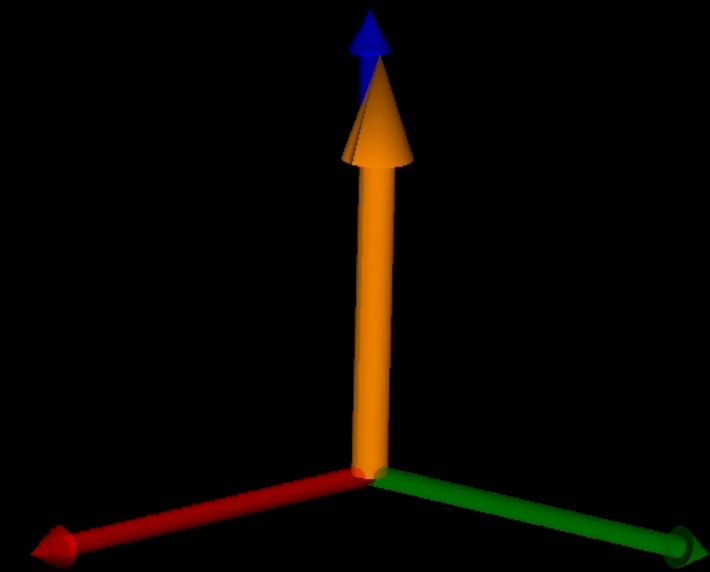
Spin Gymnastics - Lab Frame



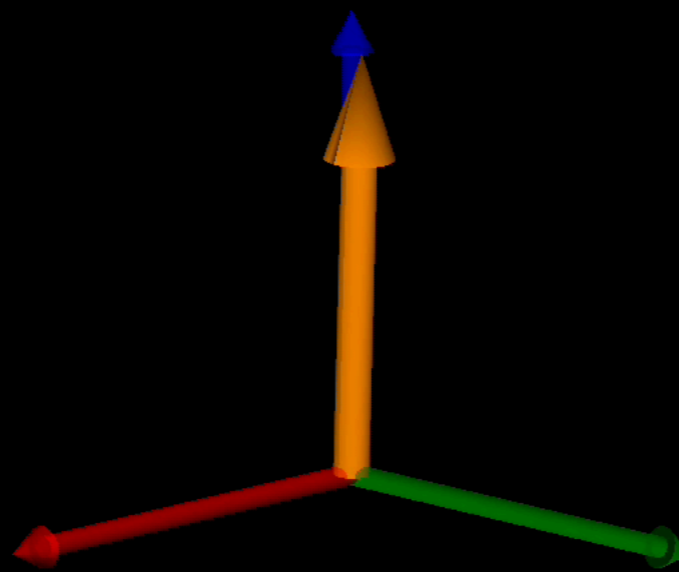
Spin Gymnastics - Rotating Frame

$$M_Z(t) = M_Z^0 e^{-\frac{t}{T_1}} + M_0 \left(1 - e^{-\frac{t}{T_1}}\right)$$

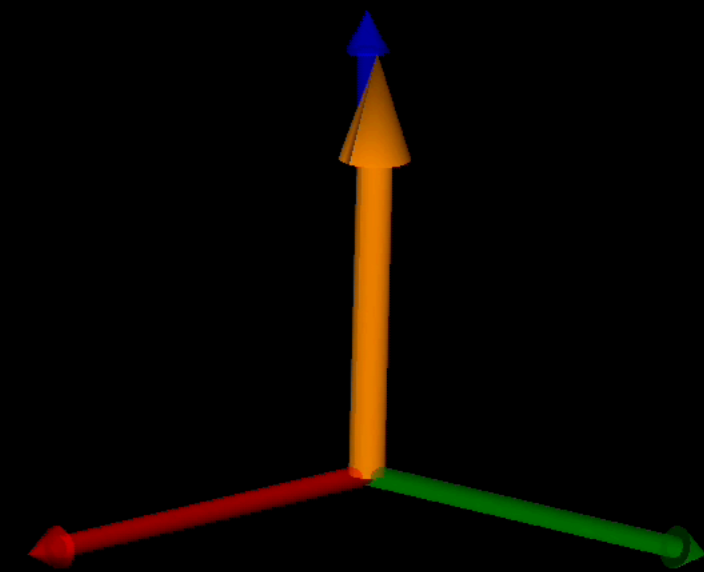
$$M_{xy}(t) = M_{xy}^0 e^{-t/T_2}$$



90° RF



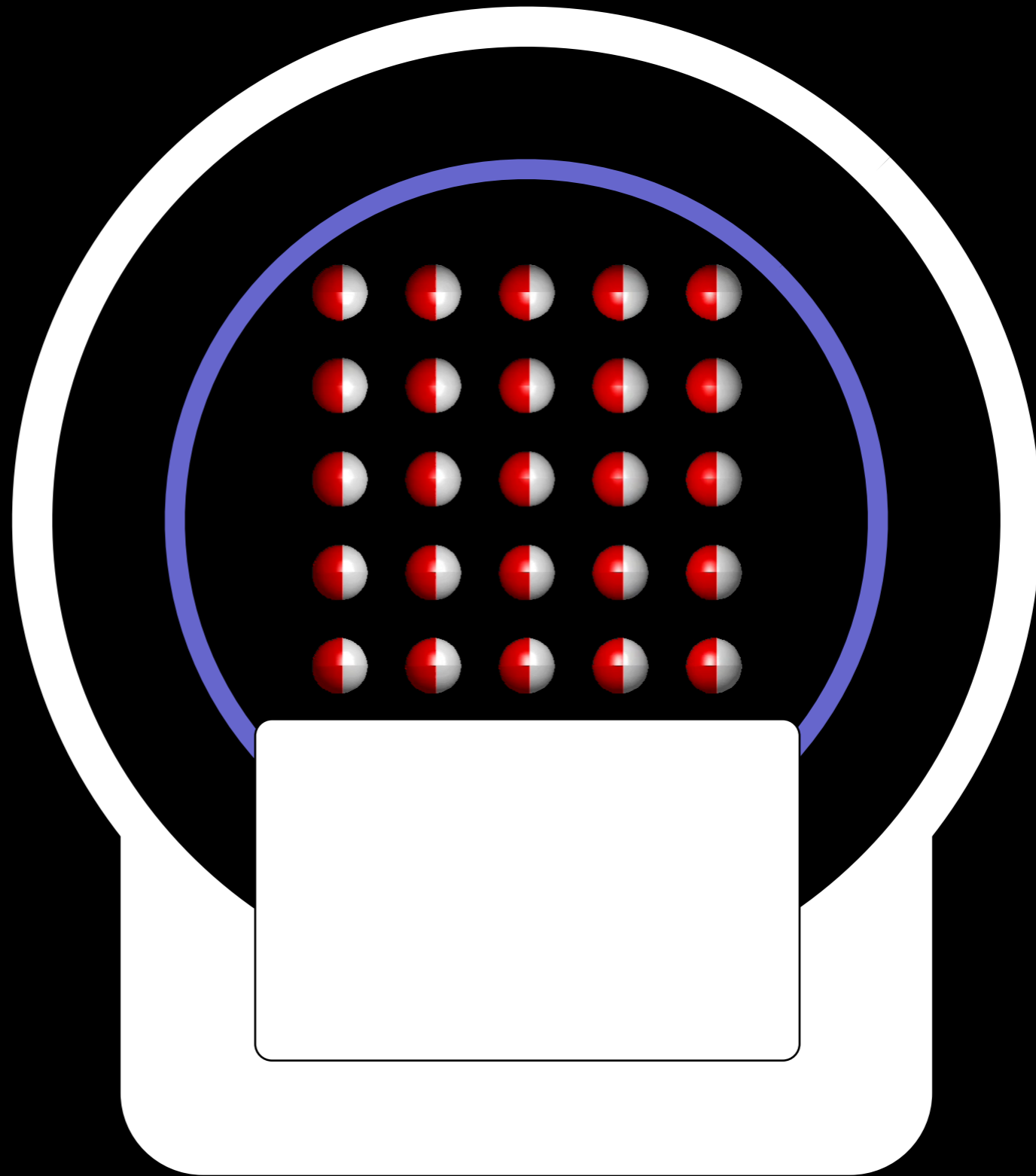
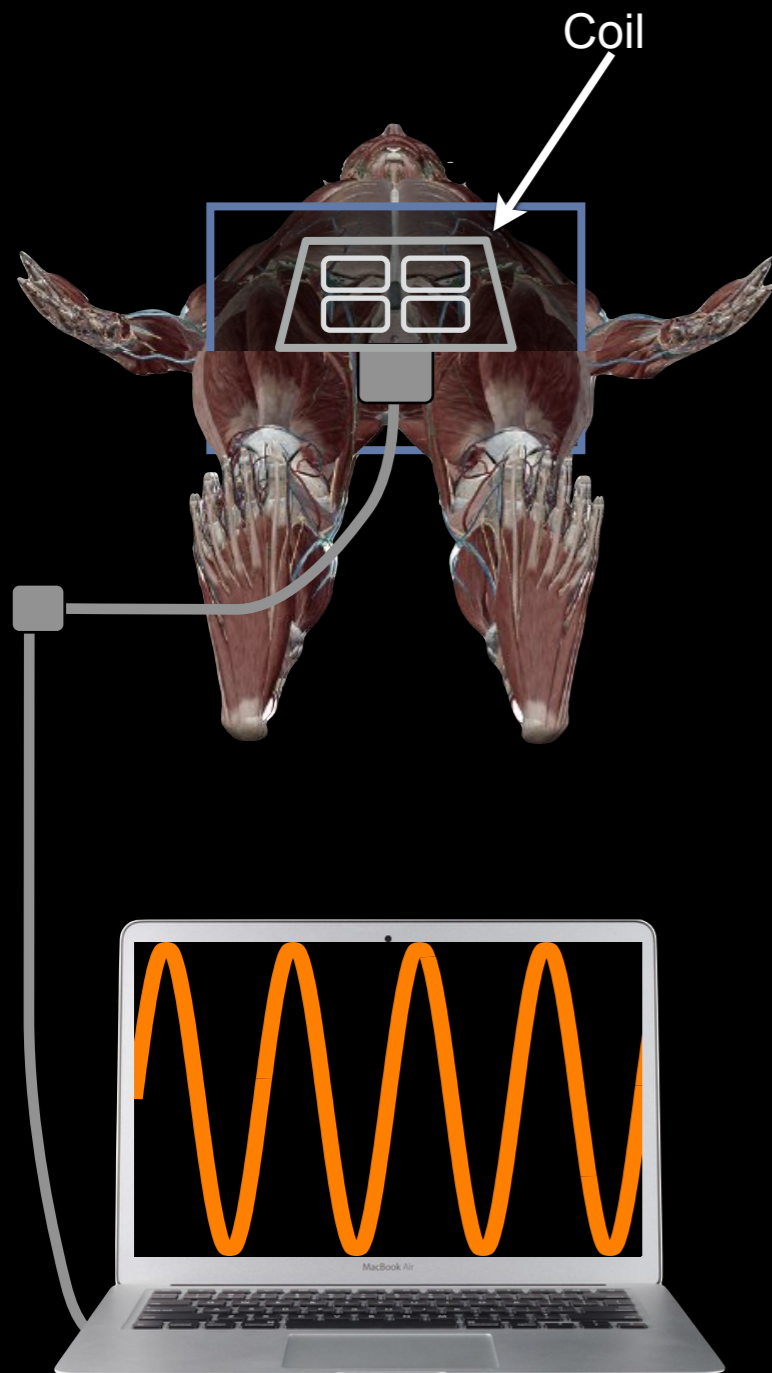
135° RF



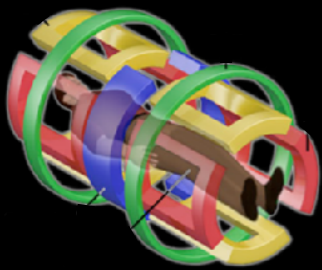
180° RF

How do we measure M_{xy} ?

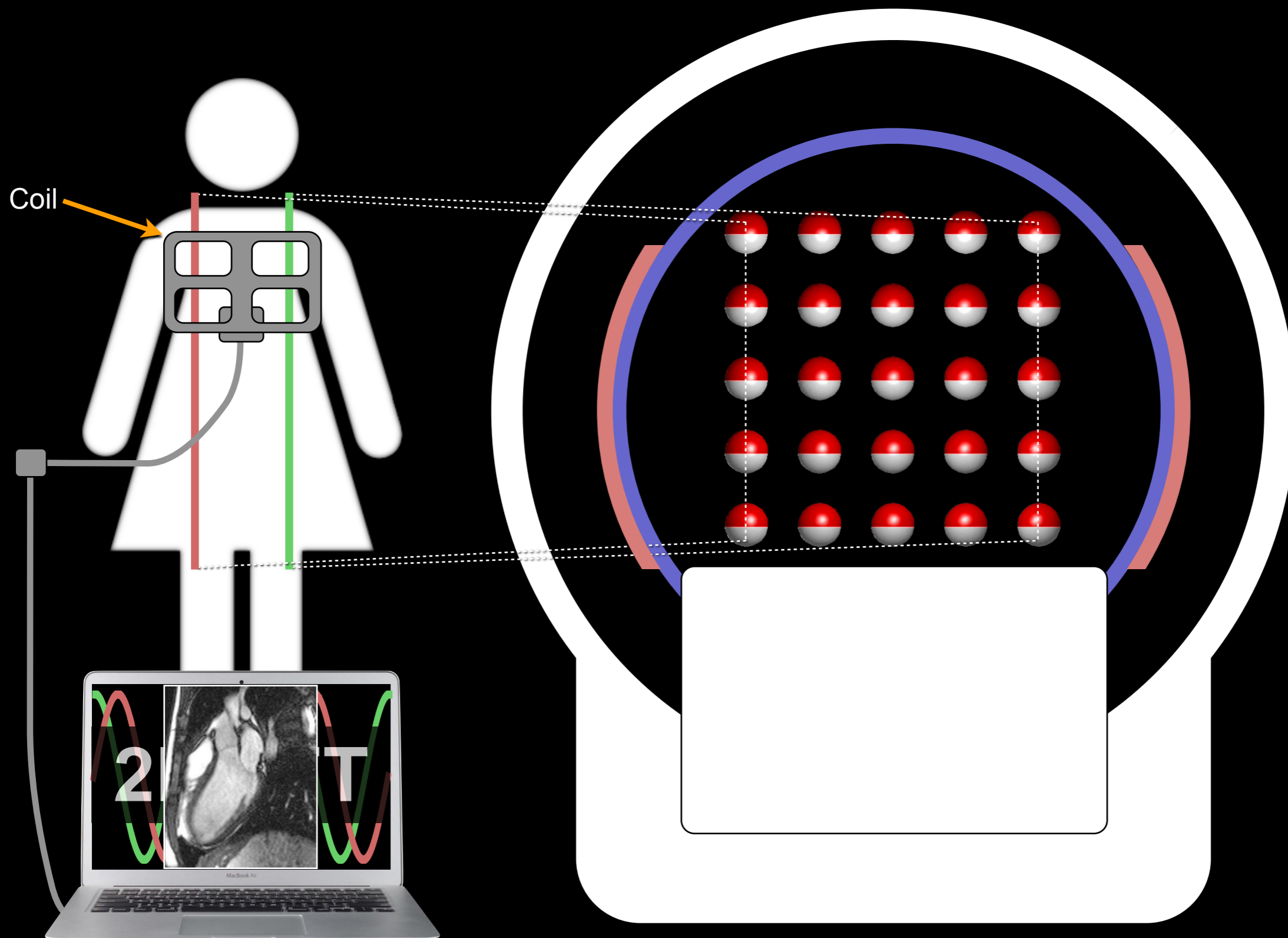
Faraday's Law of Induction



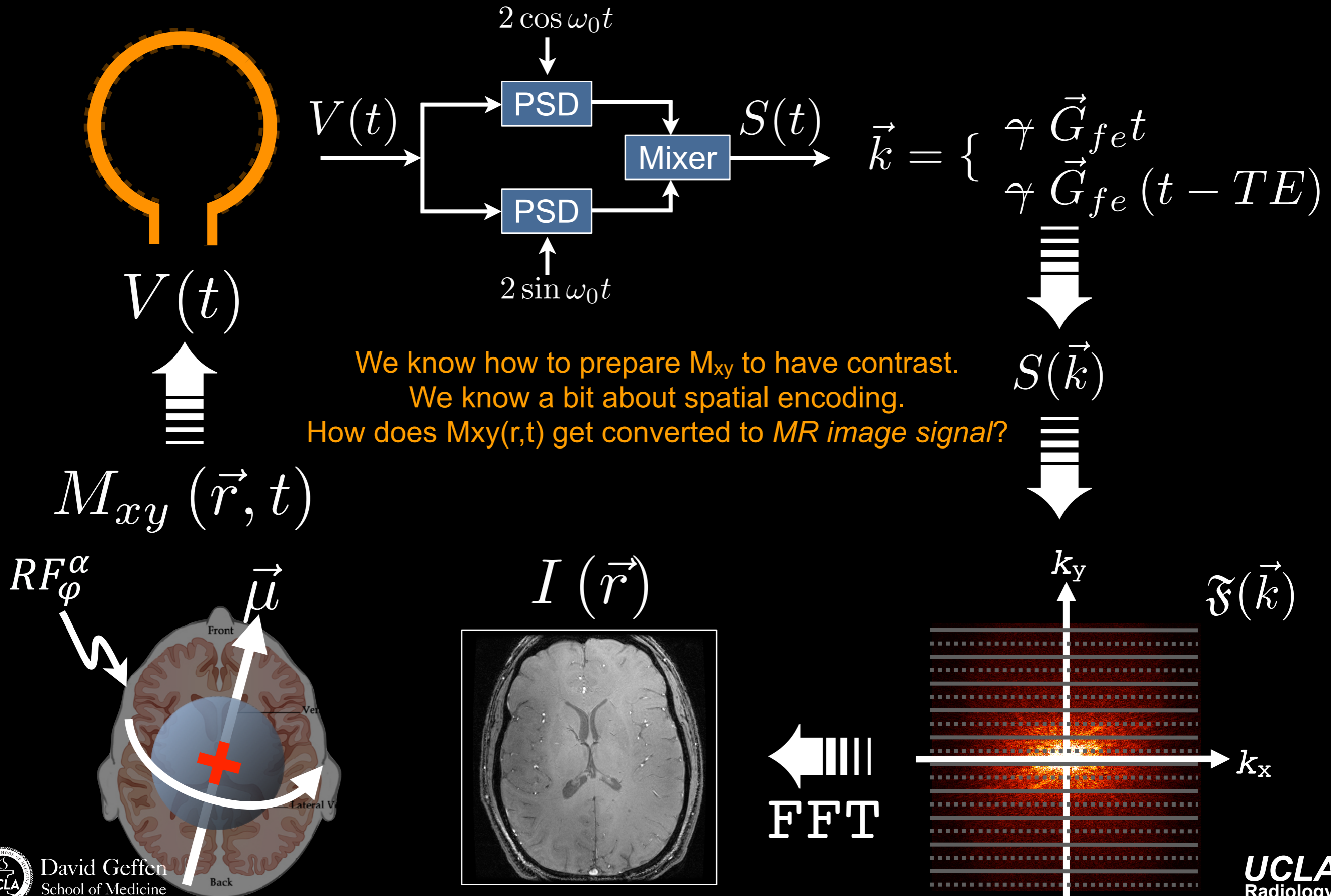
Precessing spins *induce* a current in a nearby coil.



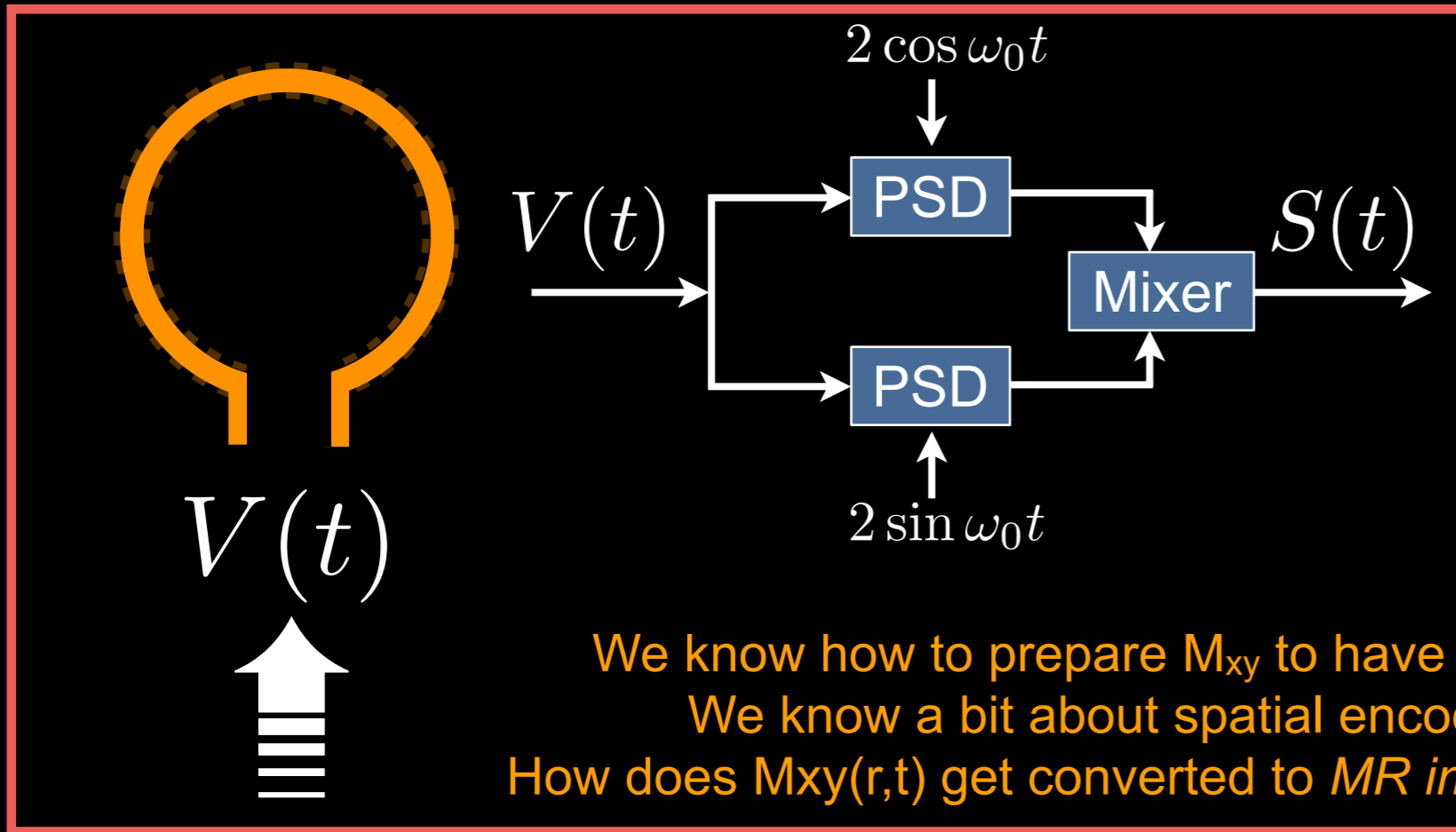
Faraday's Law of Induction



Signals in MRI

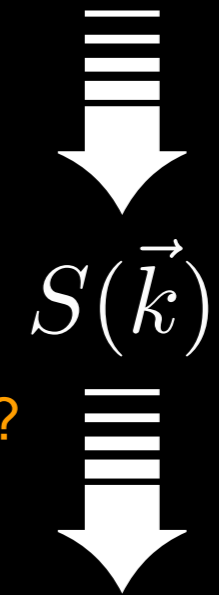


Signals in MRI



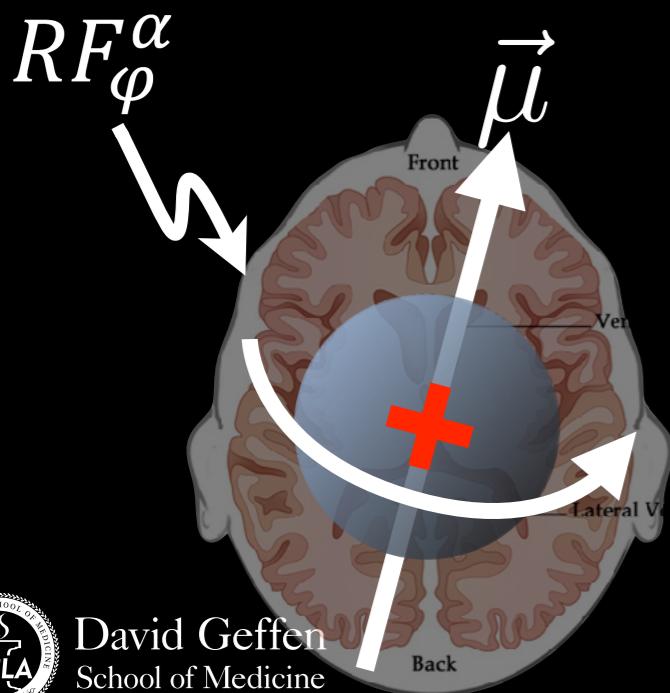
$$\vec{k} = \begin{cases} \gamma \vec{G}_{fet} \\ \gamma \vec{G}_{fe}(t - TE) \end{cases}$$

We know how to prepare M_{xy} to have contrast.
 We know a bit about spatial encoding.
 How does $M_{xy}(r,t)$ get converted to *MR image signal*?



$$S(\vec{k})$$

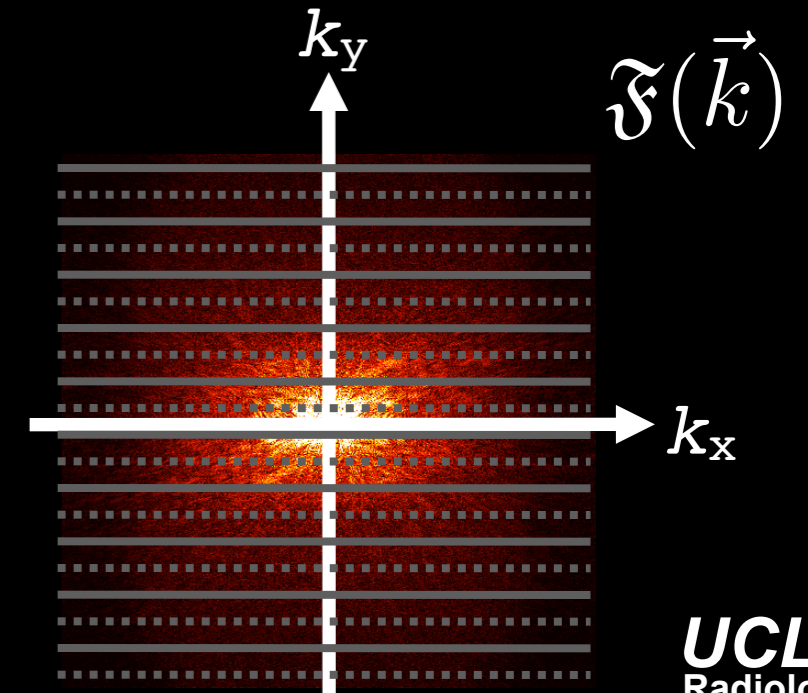
$$M_{xy}(\vec{r}, t)$$



$$I(\vec{r})$$



FFT



Basic Detection Principles

Magnetic Flux Through The Coil – *Reciprocity*

$$\Phi(t) = \int_{object} \vec{B}_r(\vec{r}) \cdot \vec{M}(\vec{r}, t) d\vec{r} \quad \text{Eqn. 3.126}$$

Magnetic Flux

Coil Sensitivity

Bulk Magnetization

The diagram illustrates the reciprocity equation for magnetic flux. The equation is $\Phi(t) = \int_{object} \vec{B}_r(\vec{r}) \cdot \vec{M}(\vec{r}, t) d\vec{r}$. Below the equation, three labels are positioned: 'Magnetic Flux' under $\Phi(t)$, 'Coil Sensitivity' under $\vec{B}_r(\vec{r})$, and 'Bulk Magnetization' under $\vec{M}(\vec{r}, t)$. Three vertical arrows point upwards from each label to its corresponding term in the equation, indicating the relationship between the physical quantities and the mathematical terms.

What happens if the coil has poor sensitivity?

What happens if the coil's sensitivity is perpendicular to the bulk magnetization? How would that happen?

Basic Detection Principles

We get here

$$S(t) = \int_{\text{object}} M_{xy}(\mathbf{r}, 0) e^{-i\gamma \Delta B(\mathbf{r})t} d\mathbf{r}$$

From Here

$$V(t) = -\frac{\partial \Phi(t)}{\partial t} = -\frac{\partial}{\partial t} \int_{\text{object}} \vec{B}(\vec{r}) \cdot \vec{M}(\vec{r}, t) d\vec{r}$$

with 25 pages of Math!

Basic Detection Principles

$$S(t) = \int_{\text{object}} M_{xy}(\mathbf{r}, 0) e^{-i\gamma\Delta B(\mathbf{r})t} d\mathbf{r}$$

Observations

Detected signal is the vector sum of all transverse magnetizations in the “rotating frame” within the imaging volume.

The Larmor frequency precession (Lab frame rotation) is necessary for detection, although only the baseband signal matters for imaging

Gradient Hardware

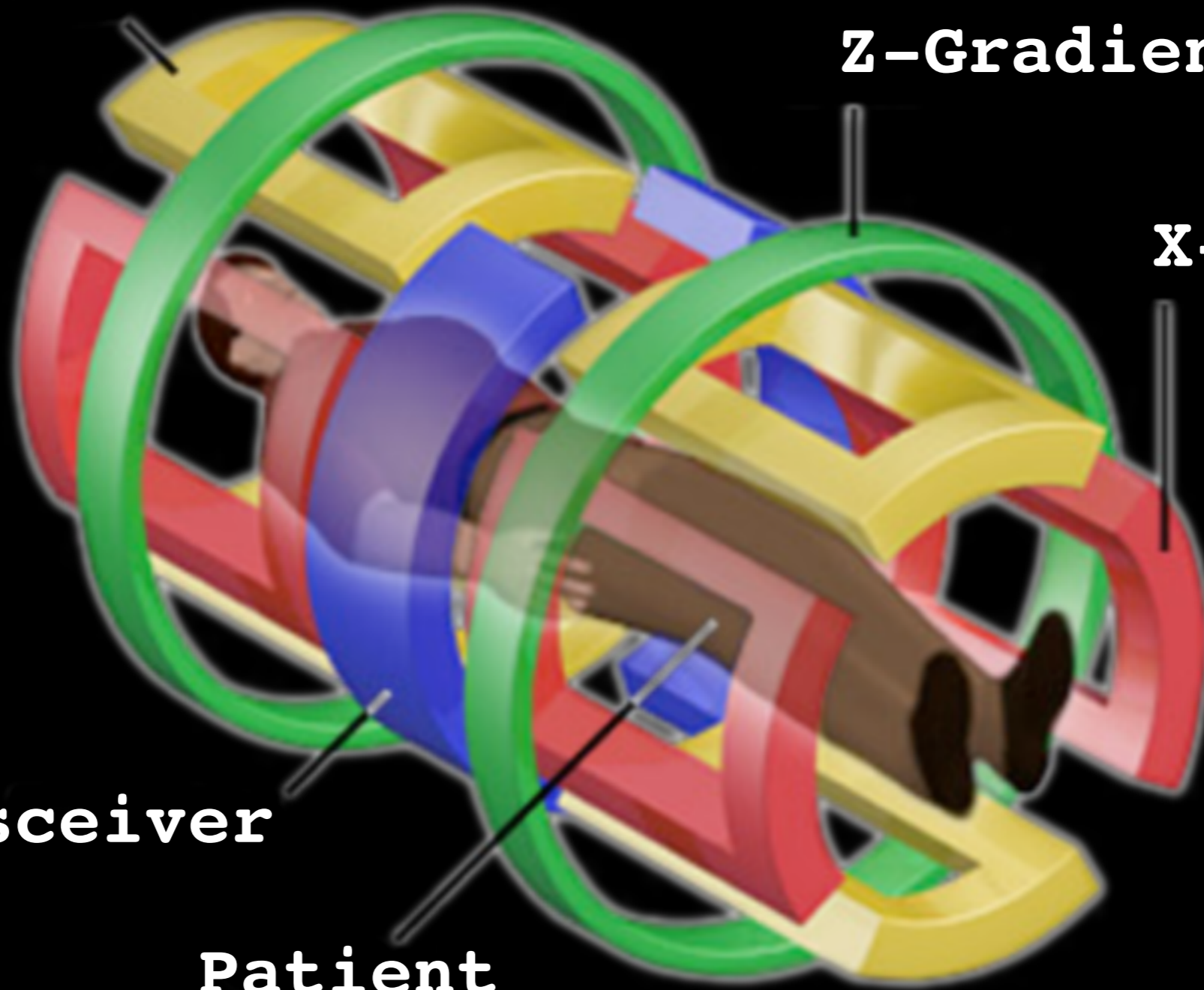
Y-Gradient

Z-Gradient

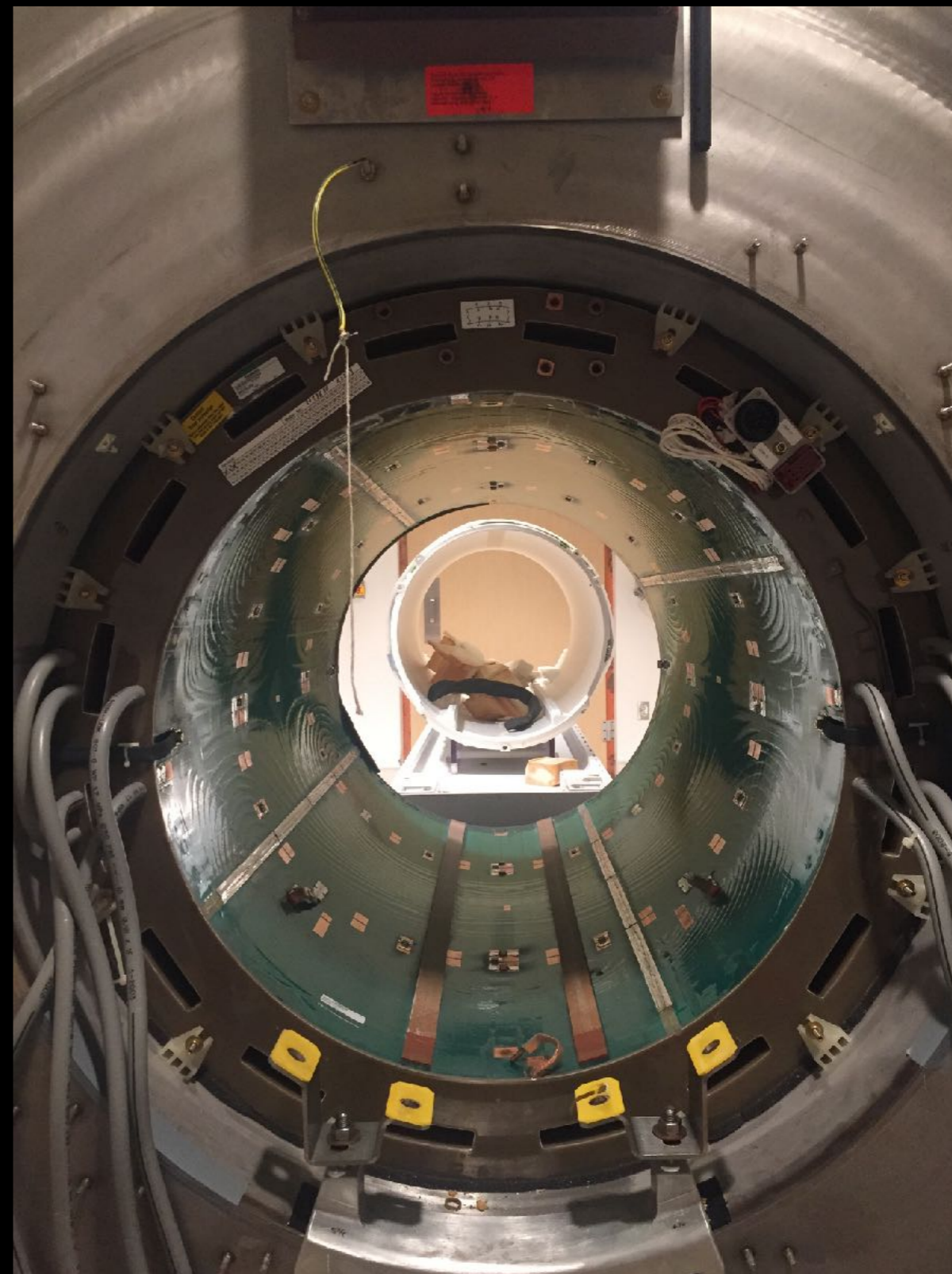
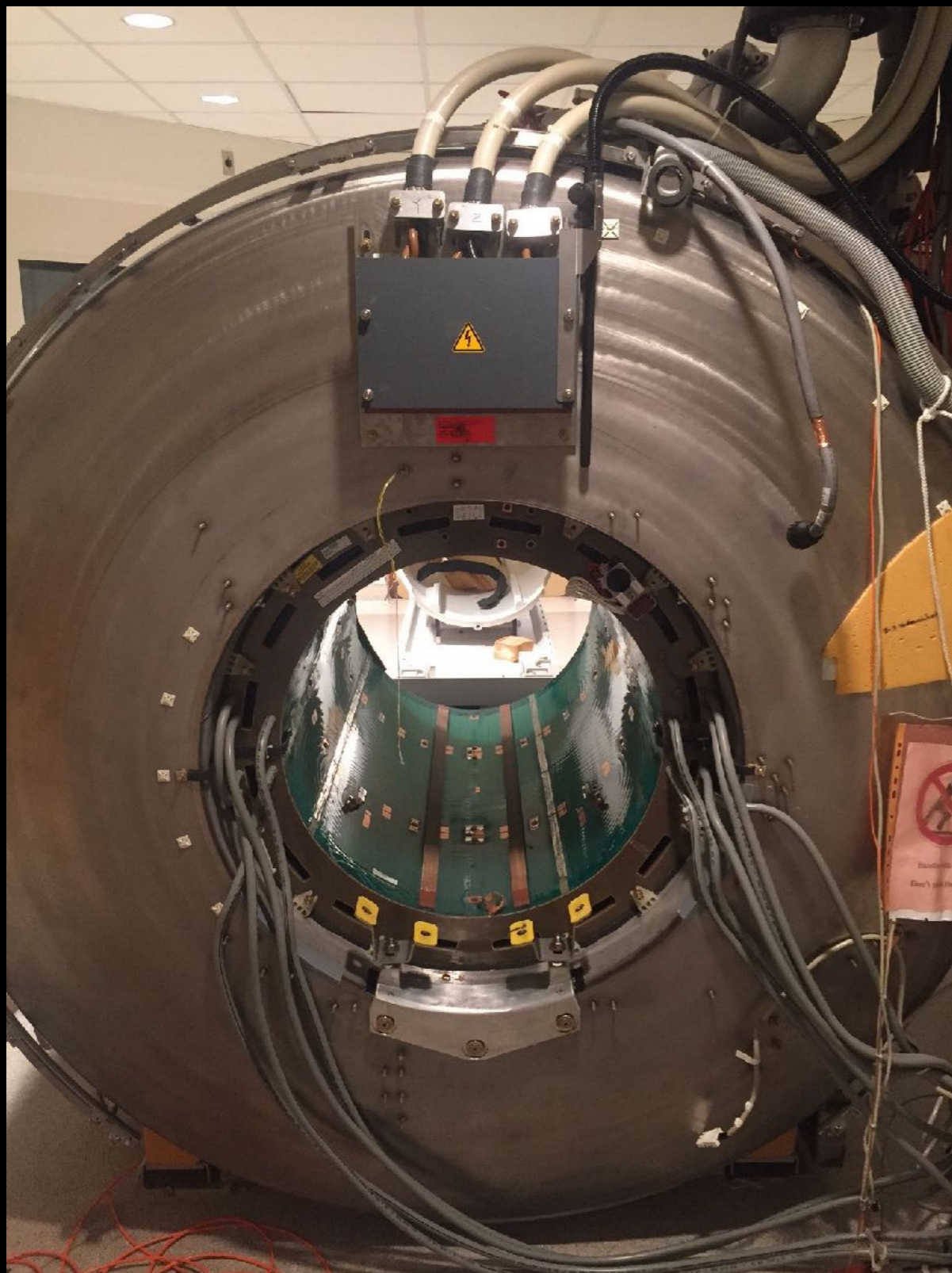
X-Gradient

Transceiver

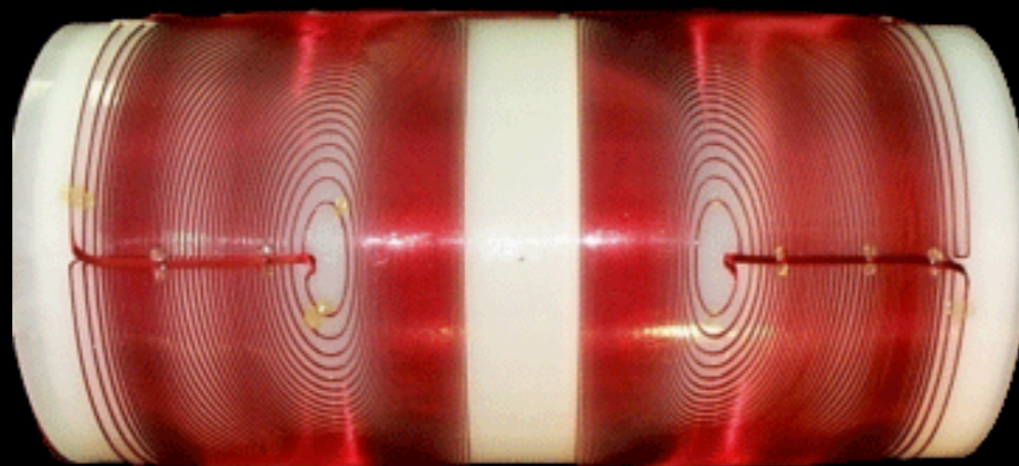
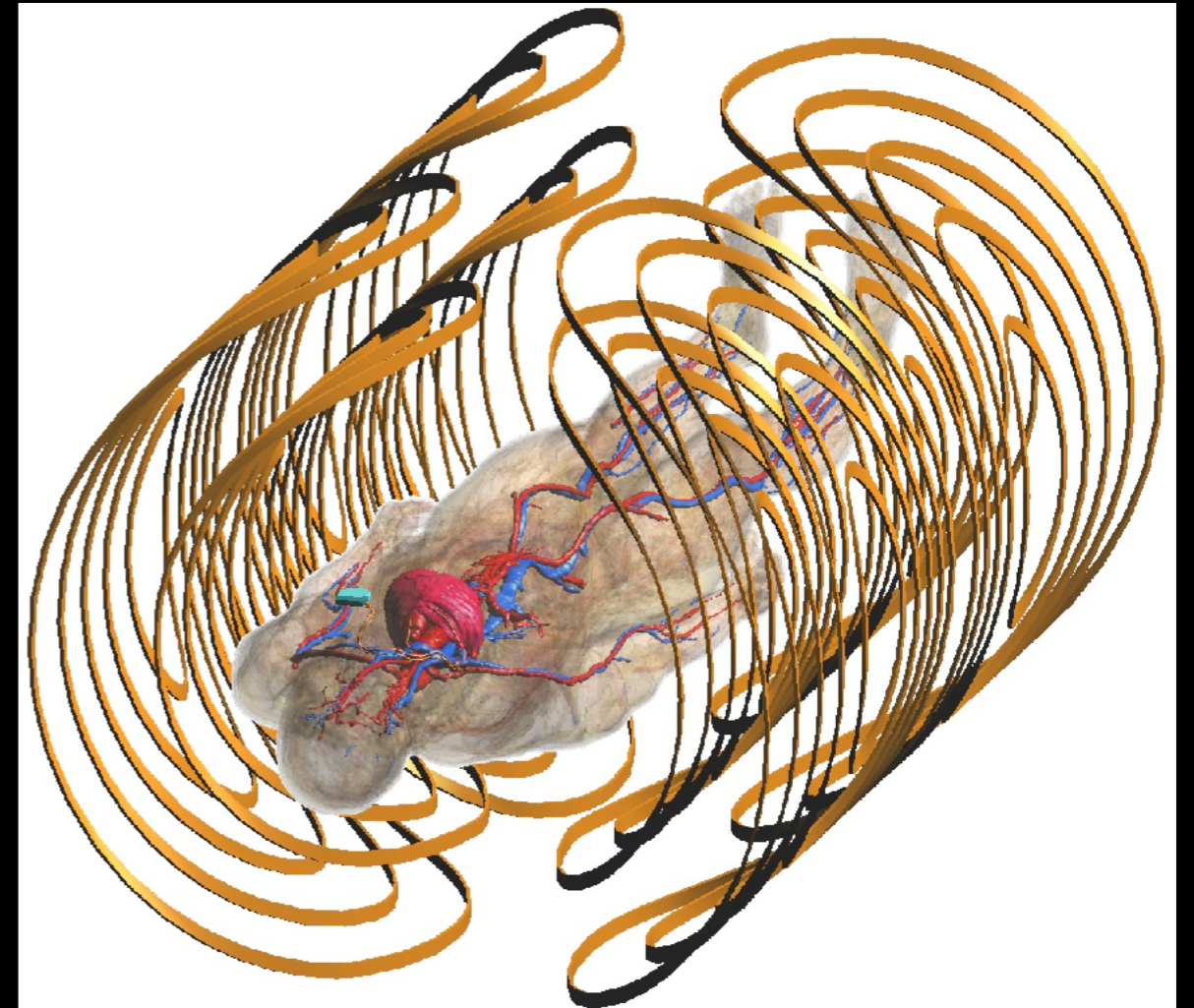
Patient



Gradient Hardware

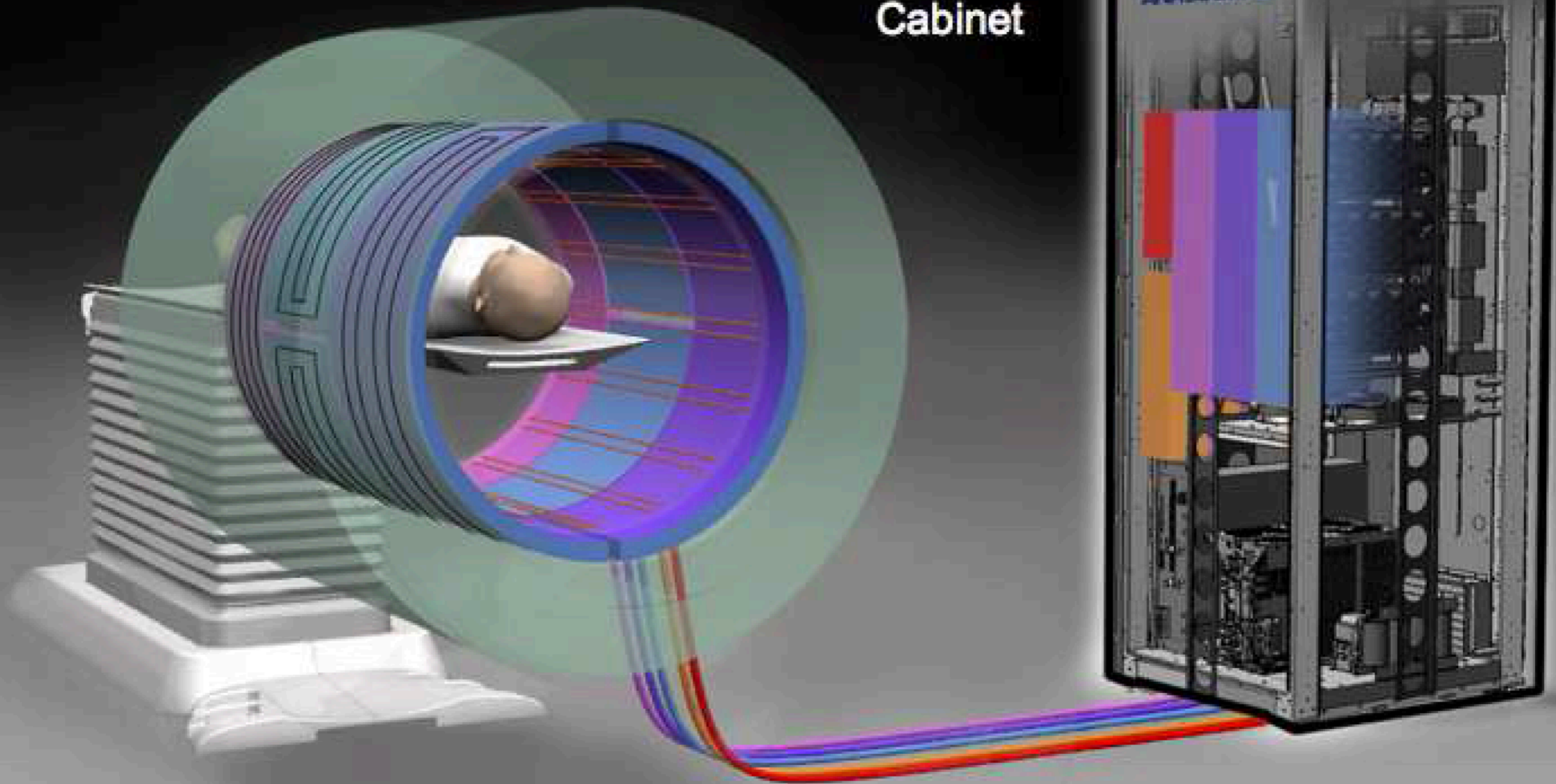


Gradient Hardware



Gradient Hardware

Integrated
MR Power
Cabinet



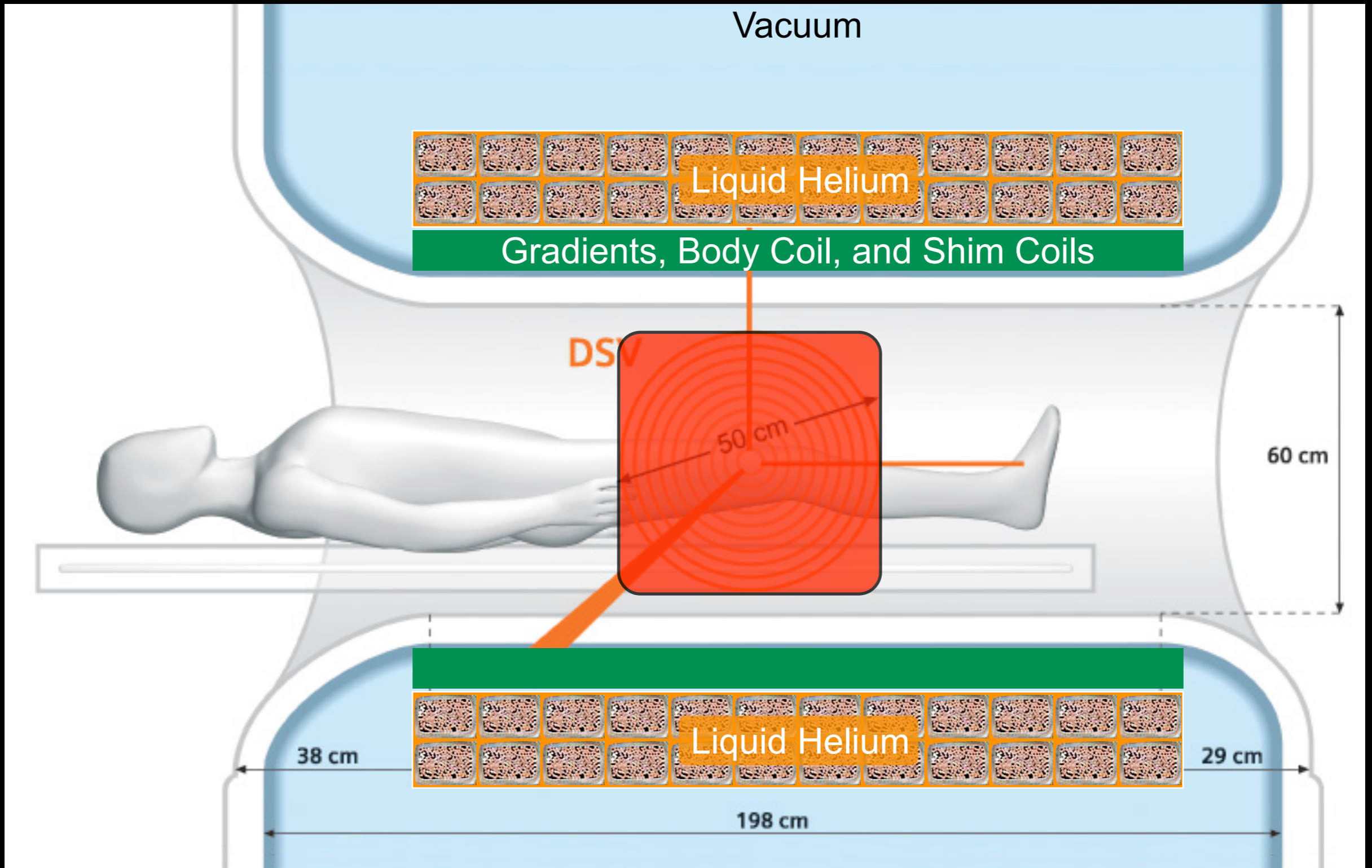
Gradients

- Primary function
 - Encode spatial information
 - Slice selection
 - Phase encoding
 - Frequency encoding
- Secondary functions
 - Sensitize/de-sensitize images to motion
 - Minimize artifacts (crushers & spoilers)
 - Magnetization **re**-phasing in slice selection
 - Magnetization **de**-phasing during readout

Gradients

- Gradients are a:
 - Small
 - $<5\text{G/cm}$ ($<0.0075\text{T}$ @ edge of 30cm FOV)
 - Spatially varying
 - Linear gradients
 - Adds to B_0 only in Z-direction
 - Time varying
 - Slewrate Max. $\sim 150\text{-}200\text{mT/m/ms}$
 - Magnetic field
 - Adds/Subtracts to the B_0 field
 - Parallel to B_0
- Gradients are NOT:
 - Fields perpendicular to B_0

Gradients



Gradients are "linear" over ~40-50cm on each axis.



Mathematics of Gradient Fields

Gradients

Gradients are a special kind of inhomogeneous field whose z-component varies linearly along a specific direction called the gradient direction.

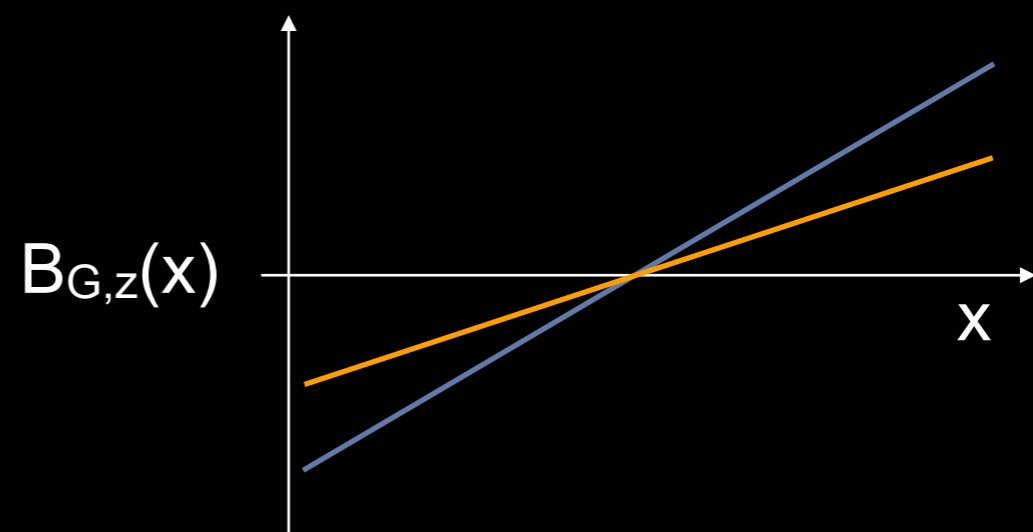
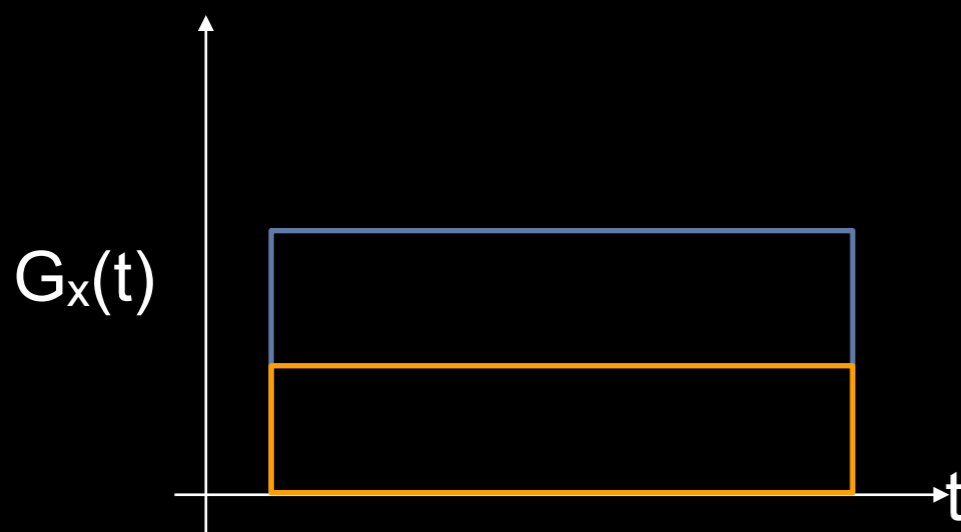
$$\underbrace{B_G}_{\text{B-field from a gradient}}, \underbrace{z}_{\text{Points along the z-direction}} \underbrace{(x)}_{\text{Varies with the x-direction}} = \underbrace{G_x}_{\text{x-gradient amplitude}} \underbrace{x}_{\text{x-distance from isocenter}}$$

Gradient Induced B-Fields

$$B_{G,z}(x) = G_x x \quad \text{x-gradient}$$

$$B_{G,z}(y) = G_y y \quad \text{y-gradient}$$

$$B_{G,z}(z) = G_z z \quad \text{z-gradient}$$



Gradient Induced B-Fields

- Each gradient coil can be activated independently and simultaneously

$$\begin{aligned} B_{G,z} \vec{k} &= (G_x x + G_y y + G_z z) \vec{k} \\ &= (\vec{G} \cdot \vec{r}) \vec{k} \end{aligned}$$

The magnetic field at a position depends on the magnitude of the applied gradient.

Combined B_0 and Gradient Fields

- Gradients contribute to the net B-field, but only along the z-direction

$$\begin{aligned}\vec{B}(\vec{r}, t) &= (B_0 + B_{G,z}) \vec{k} \\ &= \left(B_0 + \vec{G}(t) \cdot \vec{r} \right) \vec{k}\end{aligned}$$

B-Field Assumptions in MRI

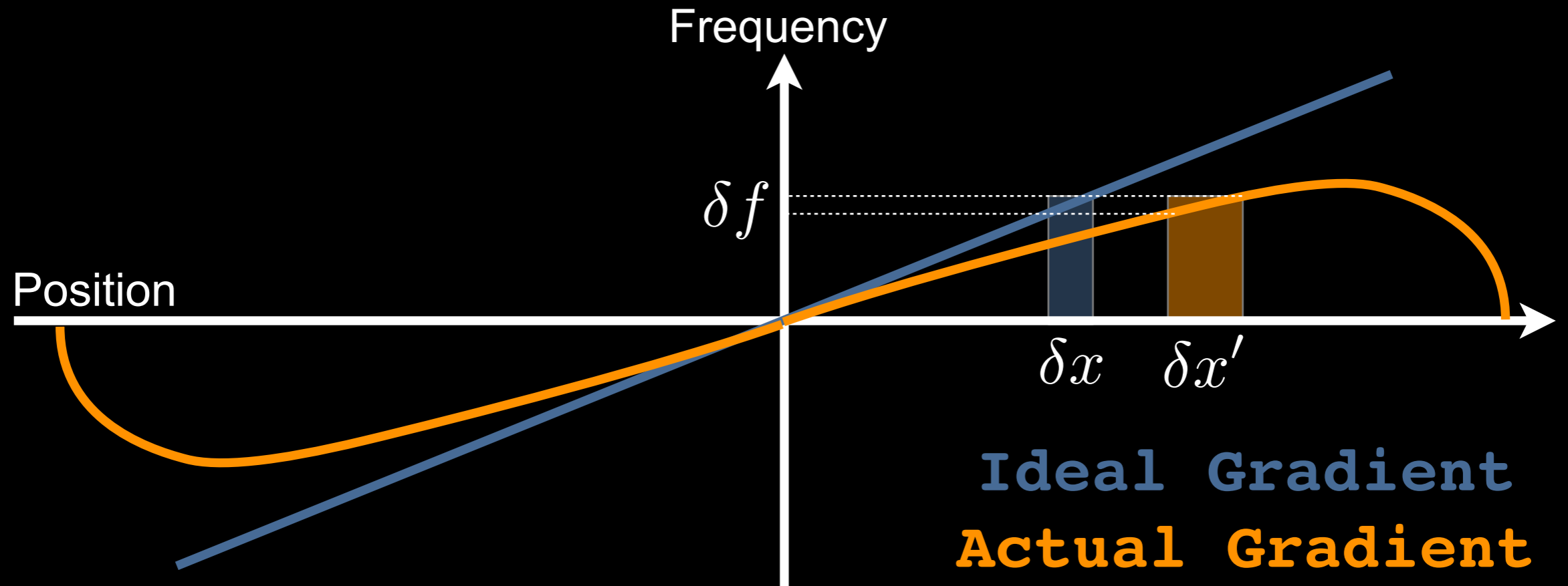
- **B_0 -field is:**
 - Perfectly uniform over space.
 - “ B_0 homogeneity”
 - Perfectly stable with time.
- **B_1 -field is:**
 - Perfectly uniform over space.
 - “ B_1 homogeneity”
 - Temporally modulated exactly as specified.
- **Gradient Fields are:**
 - Perfectly linear over space.
 - “Gradient linearity”
 - Temporally modulated exactly as specified

Imperfections of Gradient Fields

- Gradient coils aren't perfect
 - Non-linearity
 - Eddy Currents
 - Maxwell terms
(Concomitant fields)
 - But they are small
 - Much smaller than B_0

Gradient Non-linearity

Gradient Non-linearity

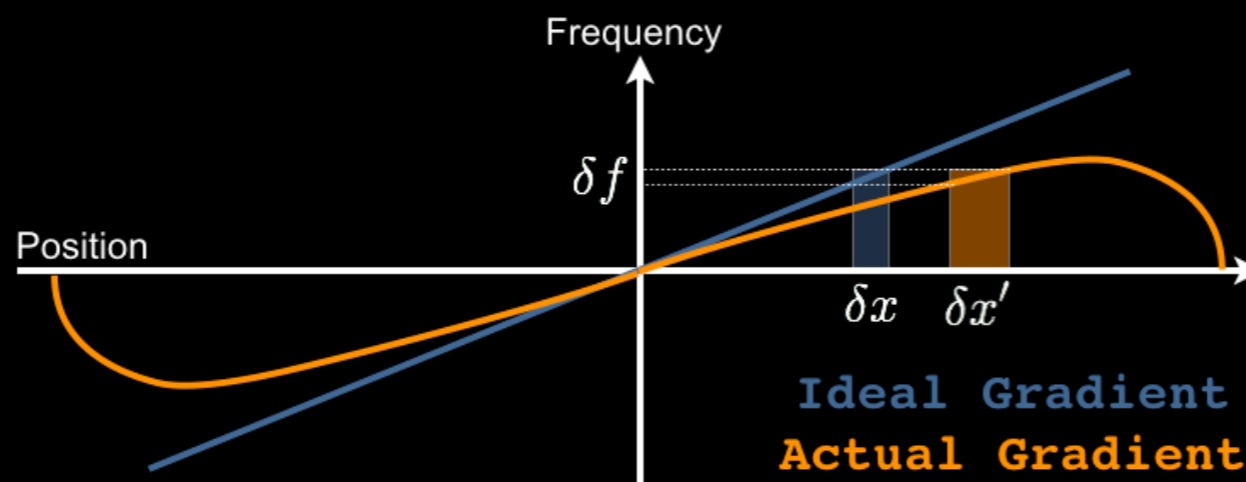


Ideally spatial position is linearly related to frequency.

Gradient Non-linearity

- Basic assumption in MRI is that the z-component of the B-field created by the gradient coils varies linearly with x, y, or z over the FOV.
- Higher gradient amplitudes and slewrates can be achieved by compromising on spatial linearity.
- Gradient non-linearity causes geometric and intensity distortions.

Gradient Non-linearity

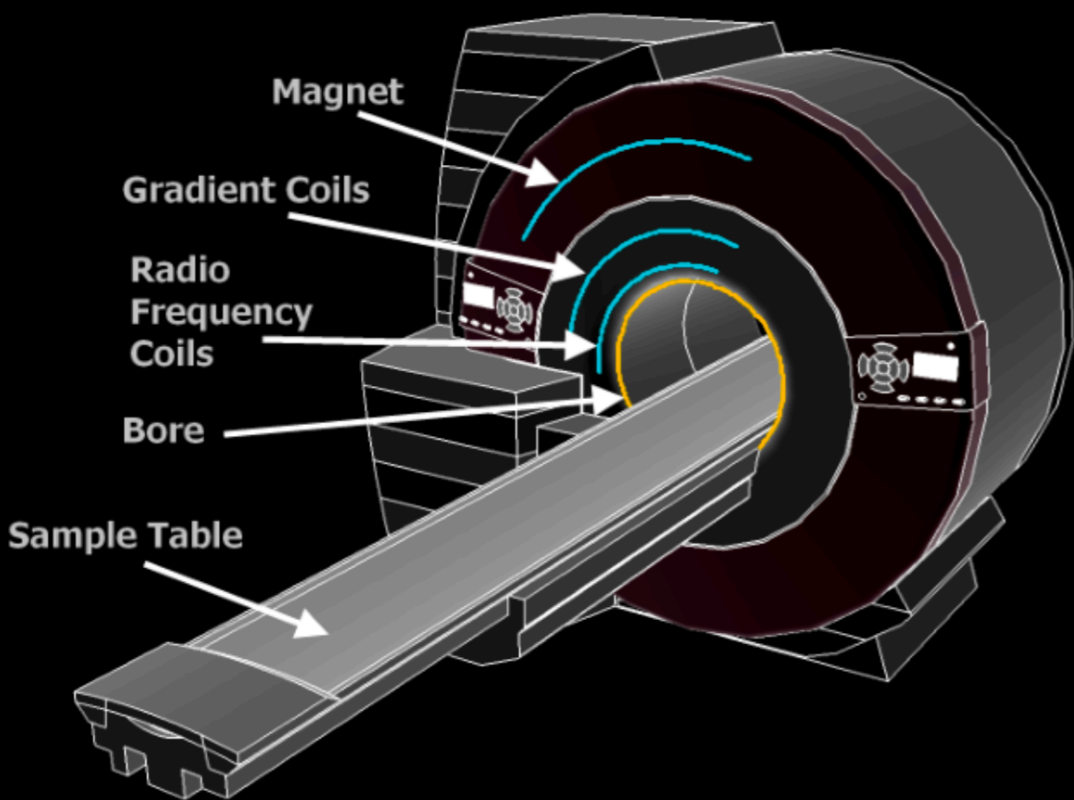


Solution

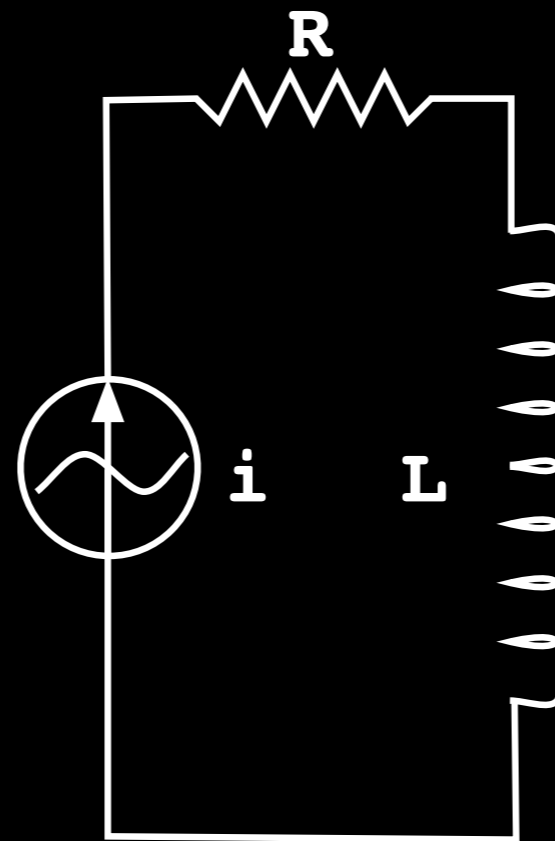
- **Improve hardware and linearity!**
- **Pay attention to FOV!**
- **Image warping parameters that are system specific and applied to all images.**
 - **Works well qualitatively.**
 - **Can be problematic quantitatively.**

Eddy Currents

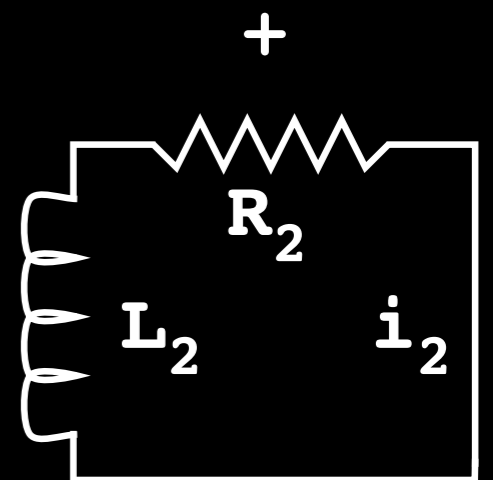
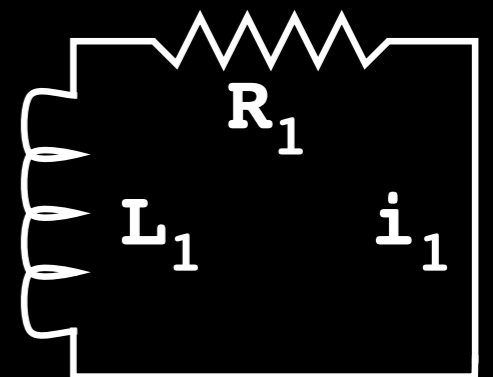
Eddy Current Origins: Diagram



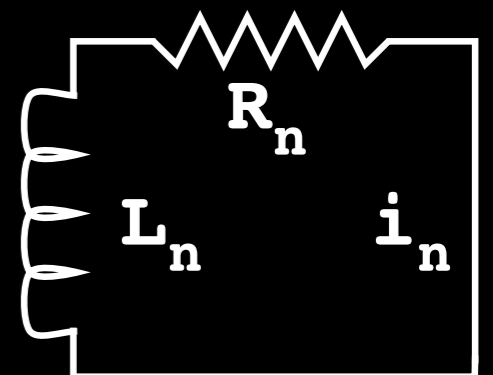
Gradient Coil



Conducting Elements



⋮

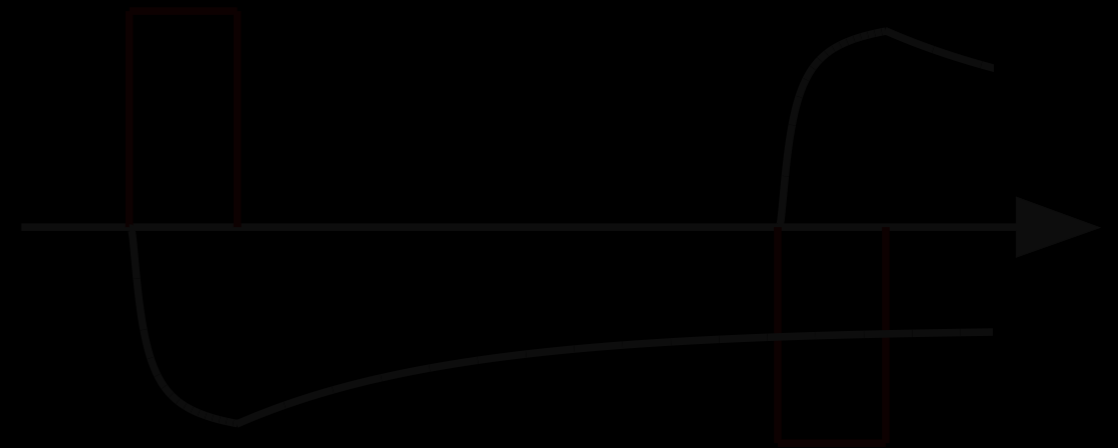


Eddy Current Gradient Distortion

Ideal Gradient Waveform



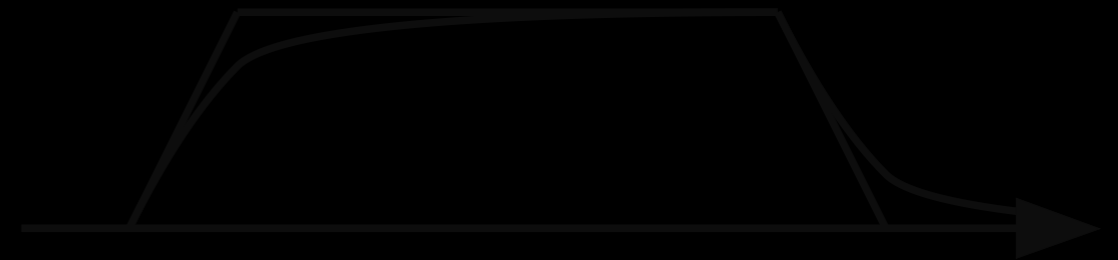
Eddy Current Gradients



Slewrate Waveform



Actual Gradient Waveform

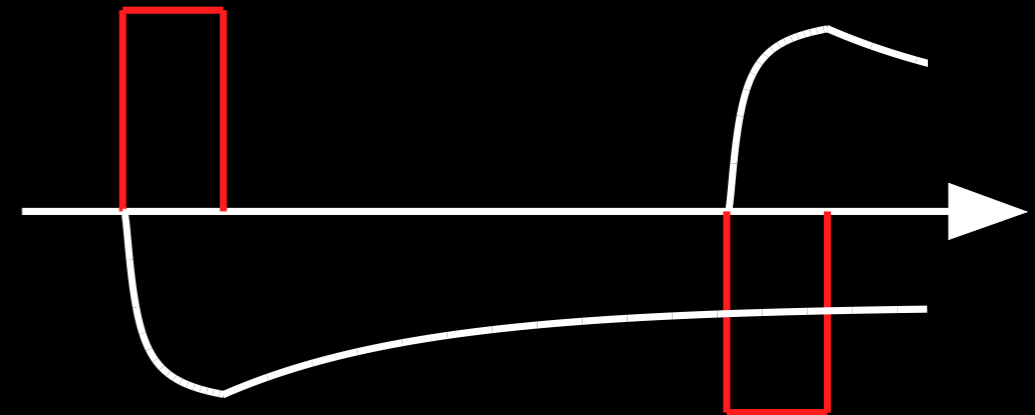


Eddy Current Origins: Mathematics

$$V_e = - \oint_{\vec{A}} \frac{\partial \vec{G}}{\partial t} \cdot d\vec{A}$$

Faraday's Law

Lenz's Law



$$I_0(t) = I_f \left(1 - e^{-\frac{Rt}{L}} \right)$$

$$I_e(t) = I_0(t_r) e^{-Rt/L}$$

$$B_e(t) \propto I_e(t)$$

Ohm's Law

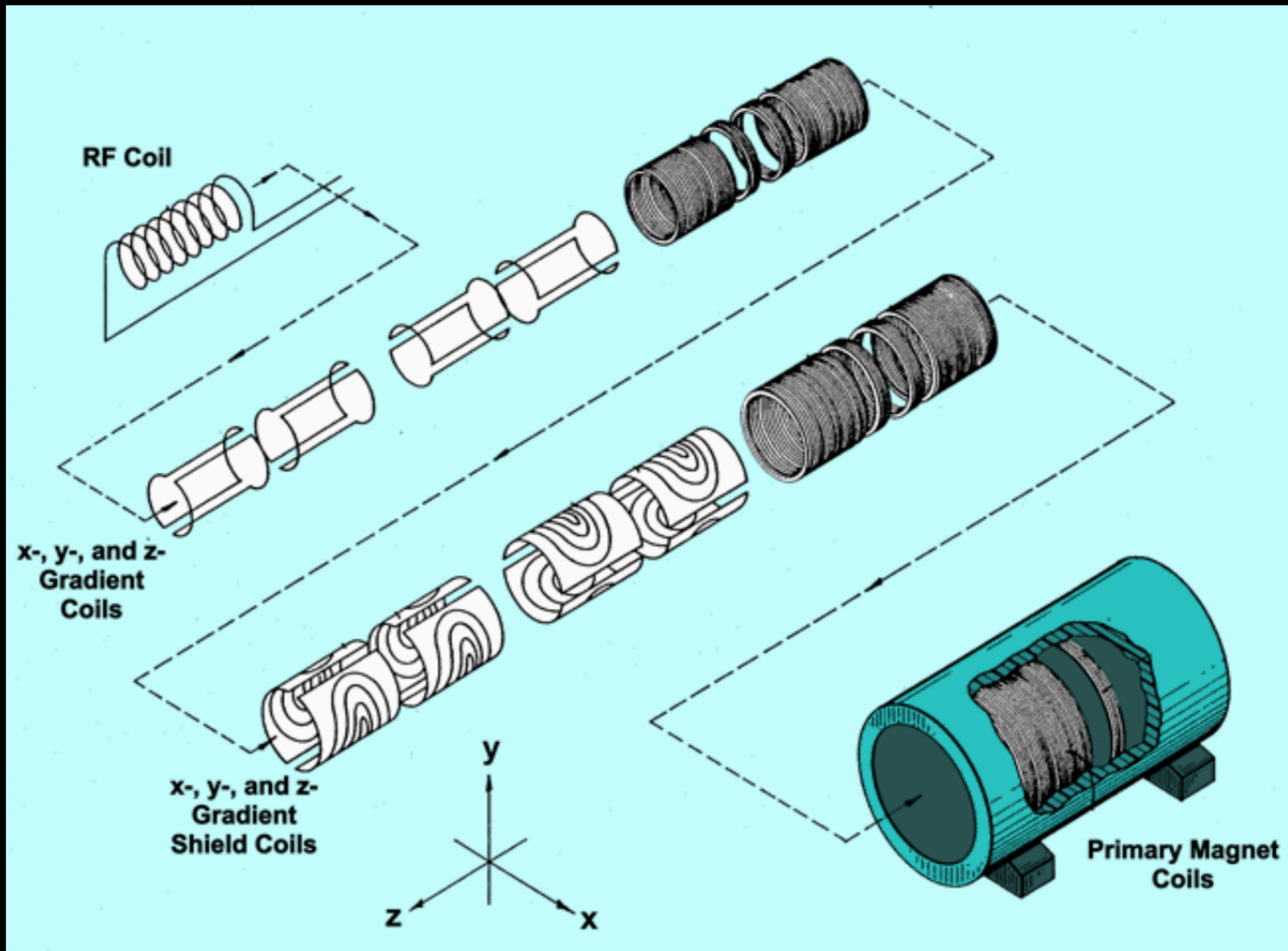
Source-Free
RL Circuit

Eddy Current
Induced Field

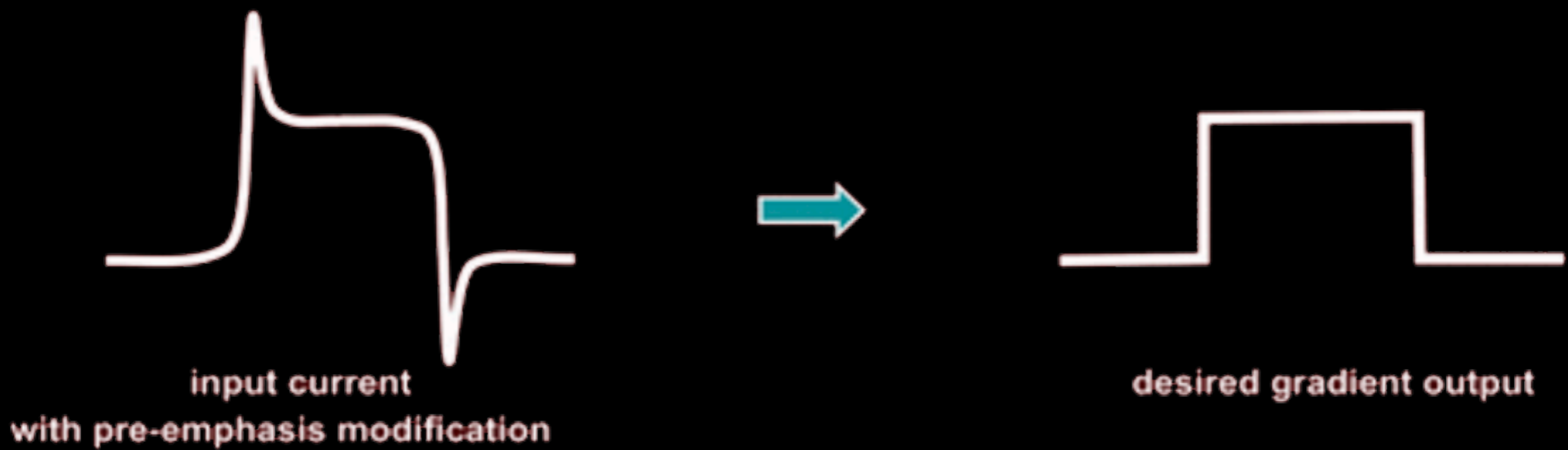
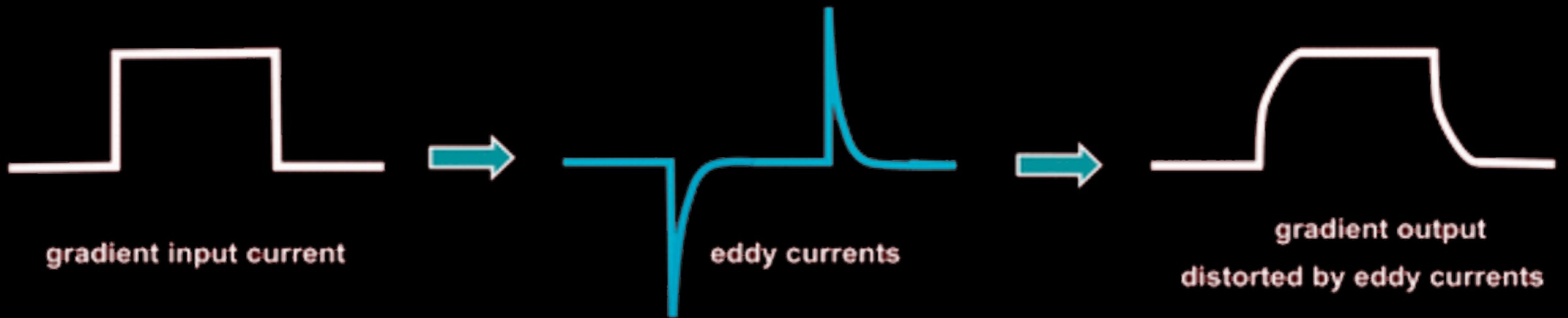
Eddy Current Compensation

- **Hardware**
 - **Actively Shielded Gradient Coils**
 - **Waveform Pre-emphasis**
- **Pulse Sequence**
 - **Slewrate de-rating**
 - **Twice Re-focused Spin Echo**
- **Reconstruction**
 - **Measure & Subtract (PC)**
 - **Predict & Subtract**

Active Shielding



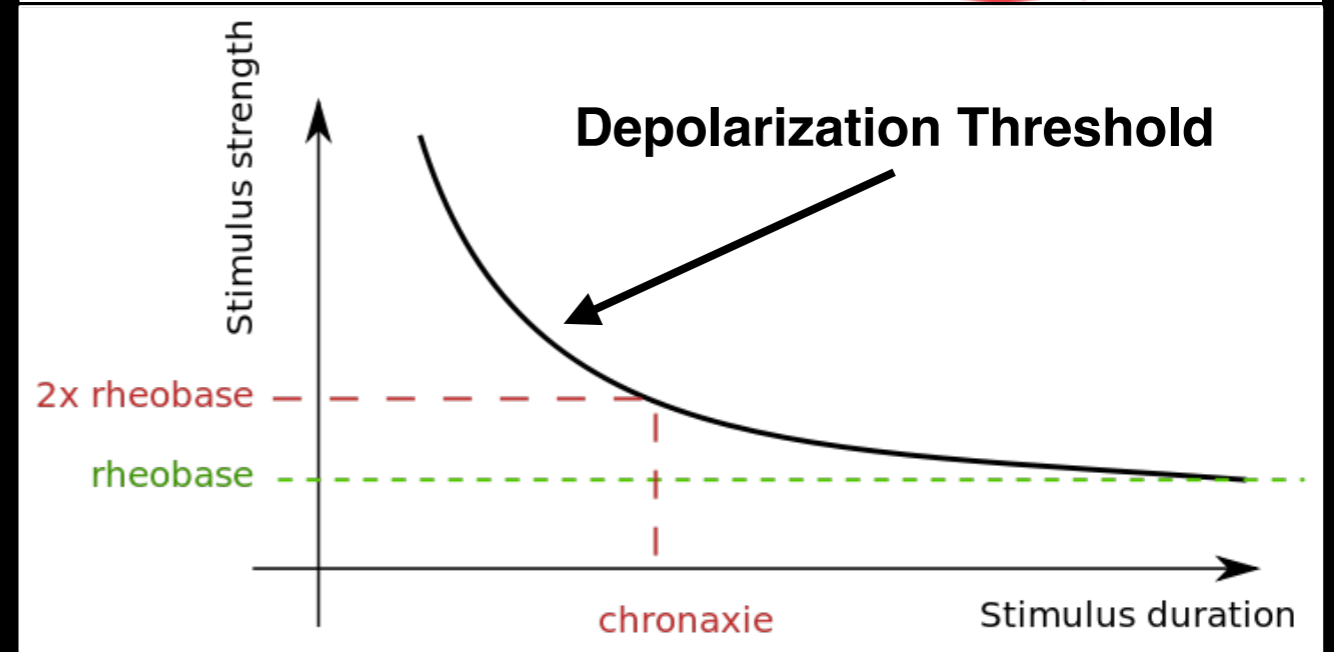
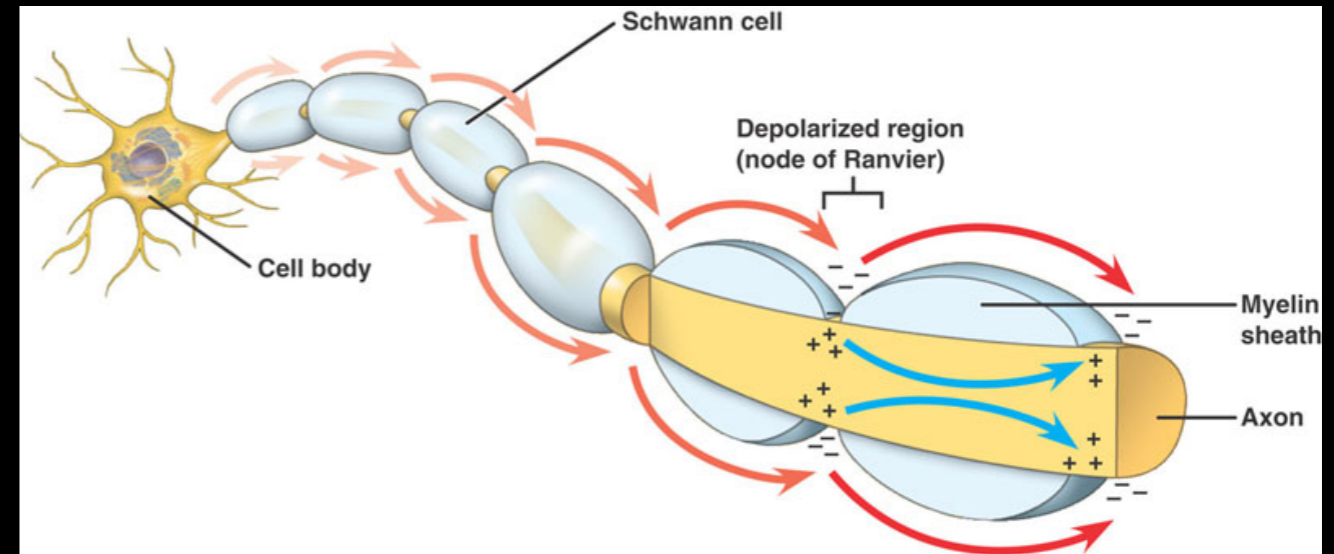
Eddy Current Pre-Emphasis



Gradient Safety

Gradient Safety

- Noise
- Peripheral nerve stimulation (PNS)



Solution: De-rate gradient slew rates, but this increases scan time.



Solution: Ear plugs



Head phones

Time-varying gradients induce mechanical vibrations and PNS.

MRI Gradient Noise



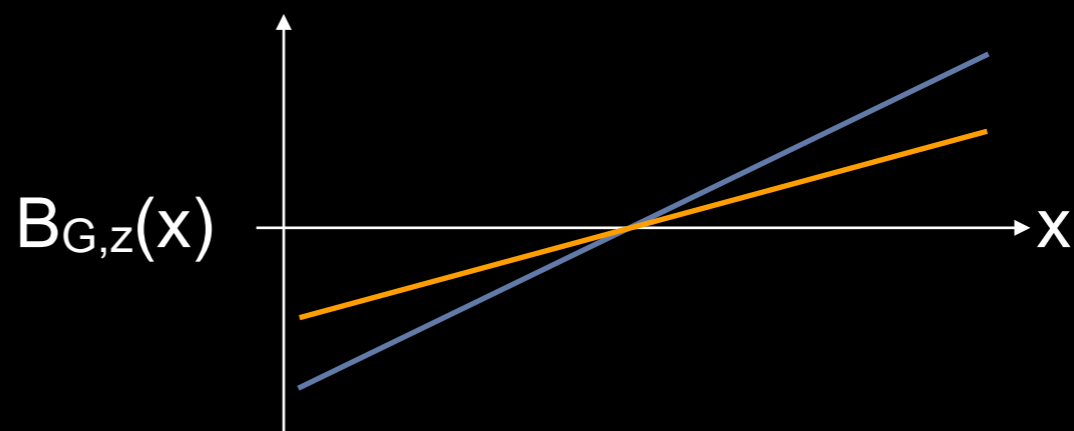
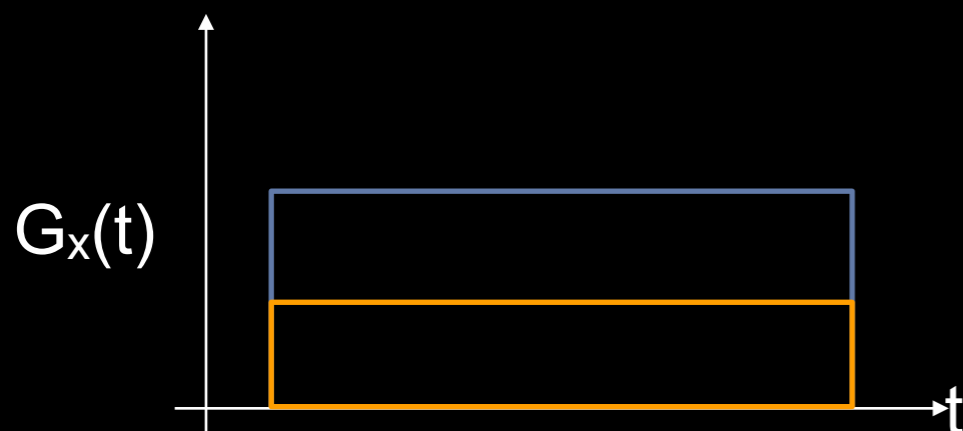
Gradient Noise

- Jet take-off @ 25m ~150 dB (eardrum rupture)
- Car horn @ 1m ~110 dB (borderline painful)
- Live rock band ~100 dB
- **MRI gradients full load** **≤99 dB**
- Garbage disposal ~80 dB
- **MRI gradients basic load** **≤75 dB**
- Radio or TV Audio ~70dB



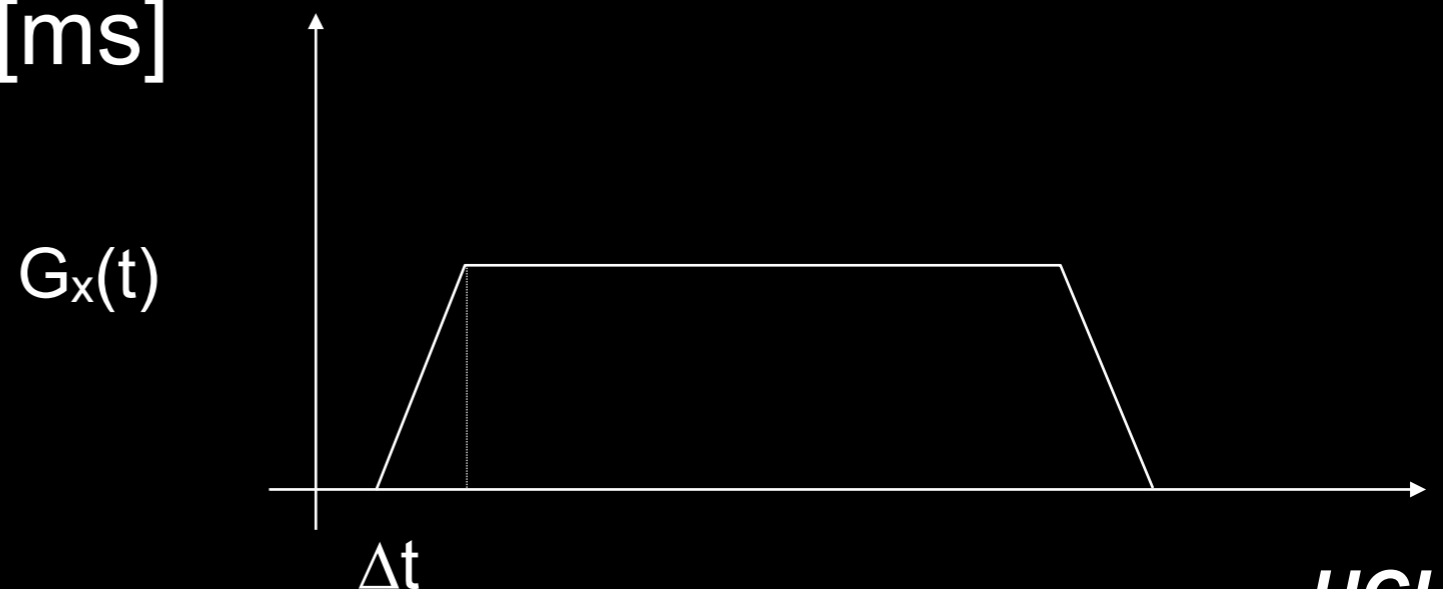
Gradient Safety – G_{Max}

- G_{max} limitations:
 - Concern: None known.
 - B_0 is already pretty big.
 - Conventional Gradients
 - $G_{Max} = 4$ to $5G/cm$ ($=50mT/m$)
 - Cutting Edge Gradients
 - $G_{Max} = 8G/cm$ ($=80mT/m$)
 - Connectome Gradients
 - $G_{Max} = 30G/cm$ ($=300mT/m$)
 - Consider the ΔB contributed by a gradient...



Gradient Slewrate

- **Gradient slew rate**
 - T/m/s (or G/cm/s)
 - dG/dt – Rate of change of gradient amplitude
- **Slew rate limited by dB/dt:**
 - Concern: Peripheral Nerve Stimulation
 - Regulated by FDA
 - Normal Mode: $dB/dt = 16 \text{ T/s} \cdot (1 + 0.36/\beta)$
 - First Level Mode: $dB/dt = 20 \text{ T/s} \cdot (1 + 0.36/\beta)$
 - $\beta = \text{stimulus duration [ms]}$



Questions?

- Related reading materials
 - Liang/Lauterbur - Chap 4.4.1
 - Nishimura - Chap 5.1.2, 7.3.2

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